

Nucleon structure studies with the PANDA experiment at FAIR

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Abstract. The PANDA experiment is one of the major projects in preparation at the upcoming FAIR facility in Darmstadt, Germany. A multipurpose high energy physics detector is currently under construction and will be operated at the High Energy Storage Ring of FAIR. High intensity antiproton beams will be available in the momentum range between 1.5 GeV/c and 15 GeV/c and will allow to address a broad physics program including hadron spectroscopy, study of charm and strangeness in nuclei, hypernuclear physics and other QCD topics. In addition, the PANDA experiment will offer unique possibilities to investigate the structure of the proton using different electromagnetic processes. In this contribution, the PANDA physics program related to the nucleon structure aspects is discussed. Feasibility studies of electromagnetic processes for the measurements of proton electromagnetic form factors in the time-like region and the nucleon-to-meson transition distribution amplitudes at PANDA, are reported.

1 Introduction

Understanding the structure of hadrons in terms of its fundamental constituents, quarks and gluons, driven by the strong interaction is one of the outstanding issues in hadron physics. The electromagnetic probes are an important tool to investigate the dynamics of Quantum Chromodynamics (QCD) in the perturbative and non perturbative regime and they continue to make important contributions to our understanding of hadron structure today. The small electromagnetic coupling of Quantum ElectroDynamics (QED) allows for a clean separation of the reaction mechanisms and the structure aspects which has led to factorization allowing for the extraction of non-perturbative structure objects of hadrons. The proton has been the subject of intensive studies since it is the lightest stable hadron. However, a complete understanding of its structure is yet to be found. For example, the proton radius, measured in muonic hydrogen atoms, has given a value that is 7 standard deviations lower than the values obtained from electronic hydrogen spectroscopy and elastic electron proton scattering [1].

Nucleon structure investigations have been performed since lepton beams became available. Elastic electron proton scattering has been intensely studied since the 1960s [2] to determine the electromagnetic form factors of the proton in the region of negative 4-momentum transfer squared ($q^2 < 0$, space-like region) [3, 4]. At high energy, one can find a hard scale such as high $Q^2 = -q^2$ so that the factorization property of the scattering amplitude can be

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37 established. Factorization theorems, derived for several classes of hard electromagnetic pro-
38 cesses, allow one to factorize the amplitudes into a perturbative hard part and a universal non-
39 perturbative hadronic matrix element containing structure functions such as the parton dis-
40 tribution functions (PDFs), generalized parton distributions (GPDs), transverse momentum
41 dependent parton distribution functions (TMD PDFs), and transition distribution amplitudes
42 (TDAs). Extensive measurements of hard (semi)-inclusive and exclusive processes have been
43 performed using high energy lepton beams [5] providing large, but not complete, amount of
44 data on these non perturbative structure functions. Lepton scattering experiments allow for
45 the access of nucleon structure matrix elements in the region of negative momentum transfer
46 q^2 of the intermediate virtual photon. Experiments accessing the region of positive, time-like
47 momentum transfer q^2 of the virtual photon, complement the studies of the nucleon structure
48 provided by the lepton beam initiated reactions. In this context, new facilities that employ
49 antiproton beams at intermediate energies like the Facility of Antiproton and Ion Research
50 (FAIR) are well suited to push back the frontiers of the accessible kinematics and achievable
51 precision. The PANDA experiment [6] at FAIR offers a unique possibility to study the proton
52 structure in time-like region with different electromagnetic processes. The physics program
53 of PANDA includes for example measurements of the time-like electromagnetic form factors,
54 the generalized distribution amplitudes (GDAs) which are the counterpart of the GPDs in the
55 time-like region, the TDAs, and the TMD PDFs. The possibility to extract the same non per-
56 turbative hadronic matrix element, like for the case of TDAs and TMD PDFs, from various
57 high energy processes and in different kinematical regimes will give experimental evidence
58 for the success of these non perturbative approaches.

59 **2 The PANDA experiment at FAIR**

60 The PANDA experiment is one of the keys projects planned at FAIR which is currently under
61 construction at Darmstadt, Germany. The PANDA experiment will measure the annihilation
62 reactions induced by a high-intensity antiproton beam with momenta from 1.5 to 15 GeV/c.
63 The PANDA detector will be installed at the High Energy Storage Ring (HESR) of FAIR
64 where first antiproton beams are expected for 2025. It is designed to provide large accep-
65 tance, high resolution and tracking capability and good neutral and charged particle iden-
66 tification in a high rate environment. The physics program includes wide range of physics
67 topics like hadron spectroscopy including charmonium and open charm states, light mesons
68 and baryons and QCD exotics; hypernuclear physics, study of hadron properties in nuclear
69 medium, nucleon structure,... In the following, the physics goals related to the measurements
70 of the proton electromagnetic form factors and the TDAs, are described.

71 **3 The Proton electromagnetic form factors**

72 Electromagnetic form factors are fundamental quantities, which describe the electric and
73 magnetic distributions of hadrons. Proton (spin 1/2 particle) is characterized by the electric
74 $G_E(q^2)$ and the magnetic $G_M(q^2)$ form factor. They are experimentally accessible through
75 measurements of differential cross sections for elastic electron proton scattering in the space-
76 like region and annihilation reactions as $e^+e^- \leftrightarrow \bar{p}p$ in the time-like region. It is assumed
77 that the interaction occurs through the exchange of one photon, which carries the momen-
78 tum transfer squared q^2 . Proton form factors have been measured in space-like region since
79 more than 60 years with increasing accuracy [2, 3]. The recent data from the JLab-GEp
80 collaboration (see [4] and references therein), based on the polarization transfer method [7]
81 in elastic electron proton scattering, showed surprisingly that the ratio $\mu_p G_E/G_M$ (μ_p is the

82 proton magnetic moment) decreases almost linearly with Q^2 . This result is in contrast to the
 83 previous measurements of unpolarized elastic ep scattering. In the time-like region, the pre-
 84 cision of the proton form factor measurements at the e^+e^- colliders and from $\bar{p}p$ annihilation
 85 experiments has been limited by the statistics. While the total cross section of the $e^+e^- \rightarrow \bar{p}p$
 86 has been measured up to high values of q^2 (~ 42 (GeV/c) 2), few data exist on the time-like
 87 proton form factor ratio in the region below $q^2 = 9$ (GeV/c) 2 . The later has been measured
 88 by PS170 at LEAR [8], BABAR [9] and more recently by BESIII [10] and CMD-3 [11]. The
 89 results of PS170 and BABAR disagree with each other with a significance up to 3σ , while
 90 the BESIII and CMD-3 measurements have large total uncertainties.

91 The PANDA experiment will allow the measurement of the proton form factors in the
 92 time-like region with the processes $\bar{p}p \rightarrow \ell^+\ell^-$, ($\ell = e, \mu$) over a large kinematical range [12,
 93 13]. It will be the first time that muons in the final state will be used to measure the time-like
 94 form factors of the proton. Moreover, the unphysical region below the kinematical threshold
 95 of the proton antiproton production of $(2M_p)^2$ can be accessed through the measurement of
 96 the $\bar{p}p \rightarrow \ell^+\ell^-\pi^0$ process [14, 15]. This region has never been experimentally accessed.
 97 The PANDA experiment can provide unique measurements of the proton form factors in this
 98 region [16]. In addition, since the time-like proton form factors are complex functions of
 99 the momentum transfer squared q^2 , an experiment with a polarized antiproton beam and/or
 100 polarized proton target would allow access to the phase difference between G_E and G_M .
 101 The feasibility of implementing a transversely polarized proton target in PANDA detector is
 102 currently under investigation.

103 3.1 Proton form factor measurements with the $\bar{p}p \rightarrow e^+e^-$ process at PANDA

104 Feasibility studies of $\bar{p}p \rightarrow e^+e^-$ for the measurements of time-like electromagnetic proton
 105 form factors at PANDA have been performed [12]. Two independent simulations have been
 106 carried out to verify the reproducibility of the results and to check systematic effects. The
 107 signal selection was performed based on the raw output of sub-detectors as well as particle
 108 identification algorithms. A suppression factor for the main background process $\bar{p}p \rightarrow \pi^+\pi^-$
 109 at the order of 10^8 has been achieved keeping a large and sufficient signal efficiency. The sys-
 110 tematic uncertainties on the extracted proton form factors have been also evaluated. With an
 111 integrated luminosity of 2 fb^{-1} , the proton form factor ratio can be determined with expected
 112 statistical (total) uncertainties of about 1% (3%) at $q^2 = 5.4$ (GeV/c) 2 up to 50% at 13.9
 113 (GeV/c) 2 (Figure 1). A separate measurement of $|G_E|$ and $|G_M|$ at PANDA is also possible.

114 3.2 Proton form factor measurements with the $\bar{p}p \rightarrow \mu^+\mu^-$ process at PANDA

115 The channel $\bar{p}p \rightarrow \mu^+\mu^-$ has been also studied at four beam momenta between 1.5 and 3.3
 116 GeV/c [13], complementing the feasibility studies of the $\bar{p}p \rightarrow e^+e^-$ channel. A method
 117 based on multivariate data classification (Boosted Decision Trees) was used to optimize the
 118 separation of the signal from the main background channel $\bar{p}p \rightarrow \pi^+\pi^-$. Signal to back-
 119 ground ratios between 1:5 and 1:13 (background rejection factor of $\sim 10^{-5}$) are achieved. A
 120 subtraction of the residual background events was applied in the analysis. After the back-
 121 ground subtraction, the proton form factors are extracted from fitting the efficiency corrected
 122 angular distributions of the $\bar{p}p \rightarrow \mu^+\mu^-$ channel. The achievable precisions of $|G_E|$, $|G_M|$ and
 123 $R = |G_E|/|G_M|$ are studied assuming an integrated luminosity of 2 fb^{-1} per beam momentum
 124 setting and including both statistical and systematical uncertainties. The results of the sim-
 125 ulations show that the proton form factors can be measured with the $\bar{p}p \rightarrow \mu^+\mu^-$ channel at
 126 PANDA with a good precision. A total relative uncertainty on the measurement of the proton
 127 form factor ratio between 5% at 1.5 GeV/c and 37% at 3.3 GeV/c is expected.

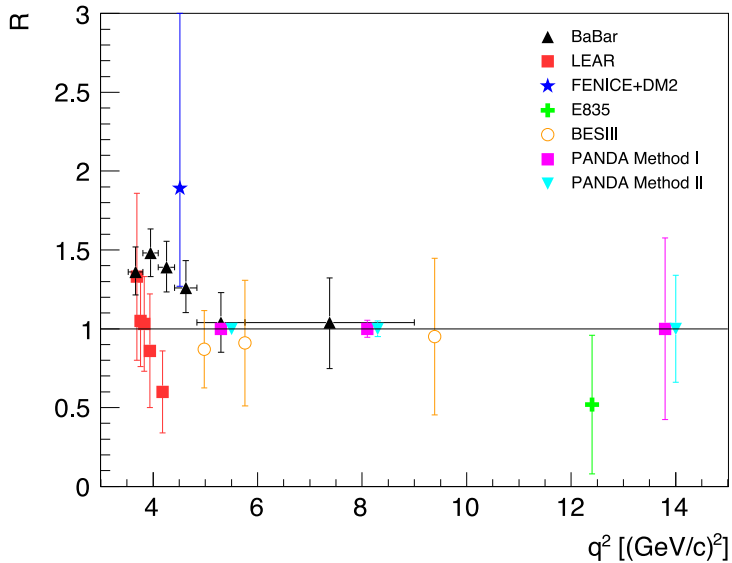


Figure 1. Expected statistical precision of the determination of the proton form factor ratio $R = |G_E|/|G_M|$ at PANDA as a function of q^2 , compared with the existing data ((see [12] and references therein)).

4 The Baryon-to-meson Transition Distribution Amplitudes

Baryon-to-meson TDAs [18, 19] were introduced as a further generalization of the GPD concept for a new class of hard electromagnetic processes in which a description, similar to the collinear factorization theorem for hard meson electroproduction [17] can be applied. Baryon-to-meson TDAs describe partonic correlations inside nucleons. Fourier transforming TDAs to the impact parameter space allows one to perform femto-photography of hadrons. In particular, nucleon-to-meson (πN) TDAs may be used as a tool for spatial imaging of the pion cloud inside the nucleon. Nucleon-to-pion TDAs arise in the description of several hard exclusive reactions such as backward electroproduction of pions off nucleons [20, 21], which can be studied at JLab [22] and COMPASS in the space-like regime. The PANDA experiment offers possibilities to access the same non perturbative functions in the time-like regime [23–25], complementing the results obtained from the measurements of lepton beam induced reactions. Recently, detailed studies of the access to πN TDAs in the reactions $\bar{p}p \rightarrow \gamma^* \pi^0 \rightarrow e^+ e^- \pi^0$ and $\bar{p}p \rightarrow J/\Psi \pi^0 \rightarrow e^+ e^- \pi^0$ in specific kinematic regimes (large q^2 and small $|t| = |(p_\pi - p_{\bar{p}})^2|$ or small $|u| = |(p_\pi - p_p)^2|$) have been presented by the PANDA collaboration in Refs. [26, 27]. The channel $\bar{p}p \rightarrow \gamma^* \pi^0 \rightarrow e^+ e^- \pi^0$ has never been measured, however some scarce data exist for $\bar{p}p \rightarrow J/\Psi \pi^0 \rightarrow e^+ e^- \pi^0$ [28, 29] which have been used to constrain the predictions for PANDA.

The feasibility of measuring the $\bar{p}p \rightarrow \gamma^* \pi^0 \rightarrow e^+ e^- \pi^0$ process with the PANDA detector has been performed at center of mass energy squared $s = 5 \text{ GeV}^2$ and $s = 10 \text{ GeV}^2$ in the kinematic regimes where the description of this process in terms of TDAs is assumed. Assuming an integrated luminosity of 2 fb^{-1} , the results shows that the differential production cross section in bins of q^2 can be measured at PANDA with an averaged statistical uncertainty of 12% at $s = 5 \text{ GeV}^2$, and 24% at $s = 10 \text{ GeV}^2$ (Figure 2).

152 The $\bar{p}p \rightarrow J/\Psi\pi^0 \rightarrow e^+e^-\pi^0$ process has been investigated at three values of the center
 153 of mass energy squared $s = 12.3 \text{ GeV}^2$, $s = 16.9 \text{ GeV}^2$ and $s=24.3 \text{ GeV}^2$. The resonant case
 154 presents the advantage of a larger cross section than the $\bar{p}p \rightarrow \gamma^*\pi^0 \rightarrow e^+e^-\pi^0$. The statistical
 155 uncertainties to measure the differential cross section as a function of the four momentum
 156 transfer in the two validity regions (small $|t|$ or small $|u|$) are expected to be about 8 – 10 %.
 157 The measurement of both resonant and non resonant channels at PANDA allows one to test
 158 the universality of πN TDAs in different kinematic ranges. In addition, the investigation of
 159 the baryon-to-photon TDAs will be possible at PANDA with the annihilation process $\bar{p}p \rightarrow$
 160 $\gamma^*\gamma \rightarrow e^+e^-\gamma$ [19]. Feasibility measurement of the $\bar{p}p \rightarrow \gamma^*\gamma \rightarrow e^+e^-\gamma$ process at PANDA
 161 is not yet performed.

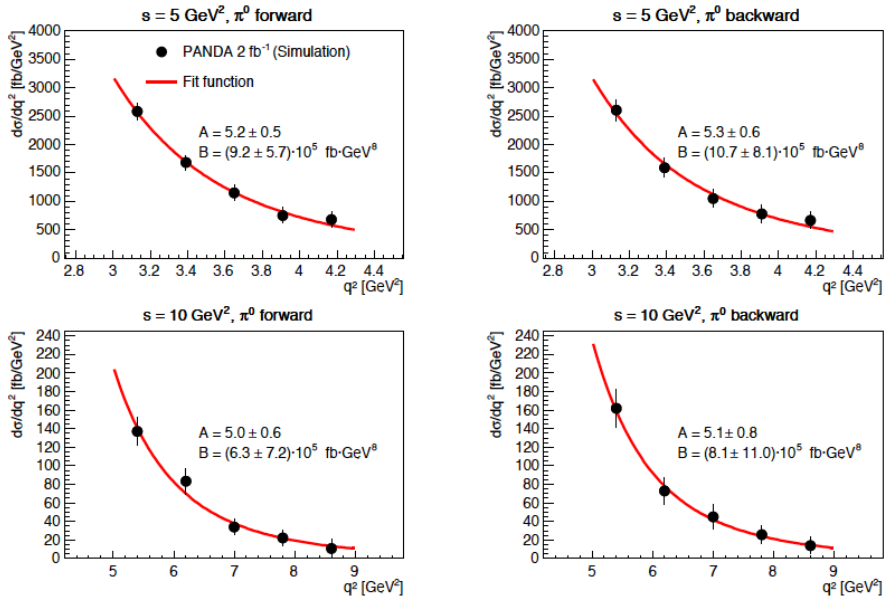


Figure 2. The $\bar{p}p \rightarrow \gamma^*\pi^0 \rightarrow e^+e^-\pi^0$ differential cross sections as a function of q^2 obtained from the simulations for $s = 5 \text{ GeV}^2$ and $s = 10 \text{ GeV}^2$, in both the t -(π^0 forward) and the u -(π^0 backward) channel kinematic regimes [26]. The resulting cross section are fitted with the theoretical leading twist predictions (see [26] and references therein). The parameters A and B are fit parameters.

162 5 Conclusion

163 The future PANDA experiment has a large and very competitive physics program to investigate the structure of the nucleon using electromagnetic processes with antiproton beams.
 164 Feasibility studies of measuring these processes with the PANDA detector have been performed, or are ongoing, using Monte Carlo simulations. In this contribution the studies related to the measurements of the time-like proton form factors and the TDAs at PANDA were presented.
 165 The results show the possibility to suppress the different background contributions to very low and sufficient levels and to extract the signal channels with good efficiencies. The PANDA experiment will allow the measurements of the proton form factors in the time-like region with high accuracy over a large and unexplored kinematical range. The resulting precisions on the measurements of the $\bar{p}p \rightarrow \gamma^*(\text{or } J/\Psi)\pi^0 \rightarrow e^+e^-\pi^0$ differential cross sections

173 will allow us in one side to test and constraint the TDA models, and on the other side to
174 measure the proton form factors in the time-like unphysical region. In addition, the precise
175 measurements of the structure functions such as the GDAs and TMD PDFs at PANDA [6]
176 and other future experiments, will provide deep insight about the internal structure of the
177 nucleon.

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