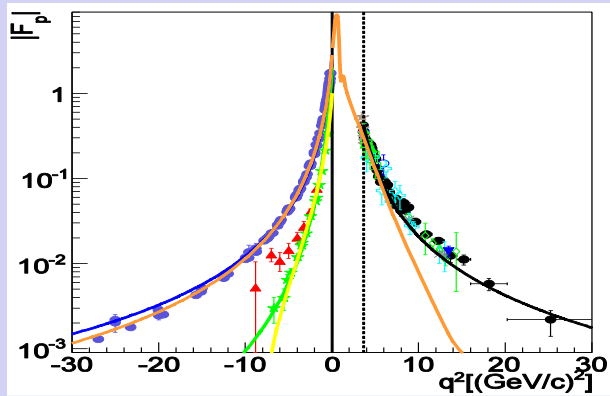


Hadron Form factors in space-like and time-like regions



Egle Tomasi-Gustafsson
IRFU, SPhN-Saclay,
and
IN2P3 - IPN Orsay France



ICNFP, July 28 - August 6, 2014

On behalf of the  Collaboration

Plan

- Introduction
 - formalism
- The Experimental Status
 - The space-like region
 - Unpolarized experiments
 - Polarized experiments
 - Issues and open questions
 - The time-like region: the PANDA Contribution
 - The unphysical region
 - The threshold region
 - The asymptotics
- Interpretation(s)
- Future prospects and Conclusions

Hadron Electromagnetic Form factors



The Nobel Prize in Physics 1961

"for his pioneering studies of electron scattering in atomic nuclei and for his thereby achieved discoveries concerning the structure of the nucleons"



Robert Hofstadter

🕒 1/2 of the prize

USA

Stanford University
Stanford, CA, USA

Characterize the **internal structure of a particle** (\neq point-like)

Elastic form factors contain information on the **hadron ground state**.

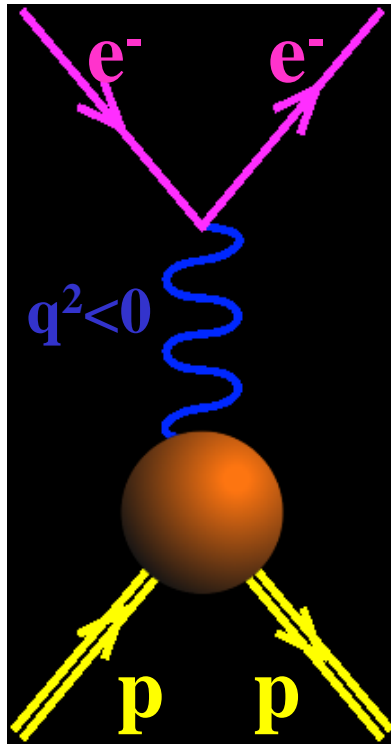
In a P- and T-invariant theory, the EM structure of a particle of spin S is defined by **$2S+1$ form factors**.

Neutron and proton form factors are different.

Deuteron: 2 structure functions, but 3 form factors.

Playground for theory and experiment at low q^2 probe **the size of the nucleus**, at high q^2 test **QCD scaling**

Electromagnetic Interaction



The electron vertex is known, γ_μ

The interaction is carried by a virtual photon of mass q^2

The proton vertex is parametrized in terms of FFs: Pauli and Dirac F_1, F_2

$$\Gamma_\mu = \gamma_\mu F_1(q^2) + \frac{i\sigma_{\mu\nu} q^\nu}{2M} F_2(q^2)$$

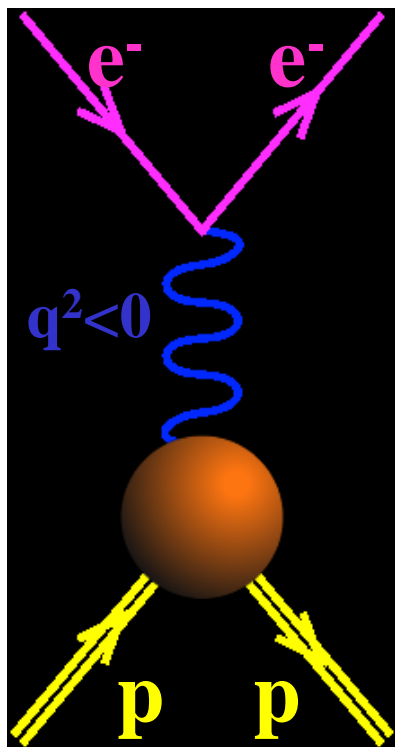
or in terms of Sachs FFs:

$$GE = F_1 - \tau F_2, \quad GM = F_1 + F_2, \quad \tau = -q^2/4M^2$$

What about high order radiative corrections?

Hadron Electromagnetic Form factors

$$\Gamma_{\mu} = \gamma_{\mu} F_1(q^2) + \frac{i \sigma_{\mu\nu} q^{\nu}}{2M} F_2(q^2)$$

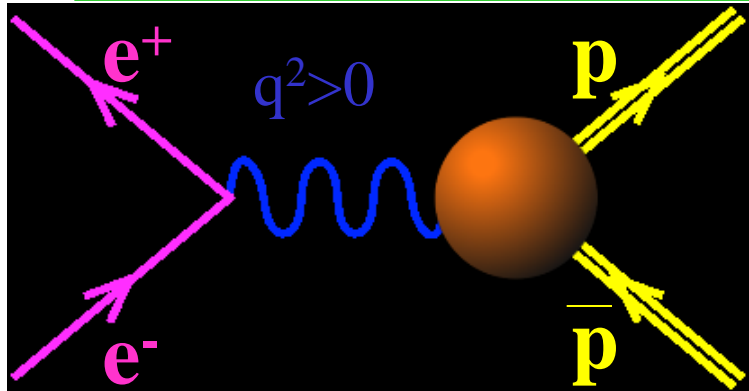


$GE(0)=1$
 $GM(0)=\mu_p$

*Space-like
FFs are real*

*Unphysical region
 $p+\bar{p} \leftrightarrow e^+ + e^- + \pi^0$*

*Asymptotics
 - QCD
 - analyticity*



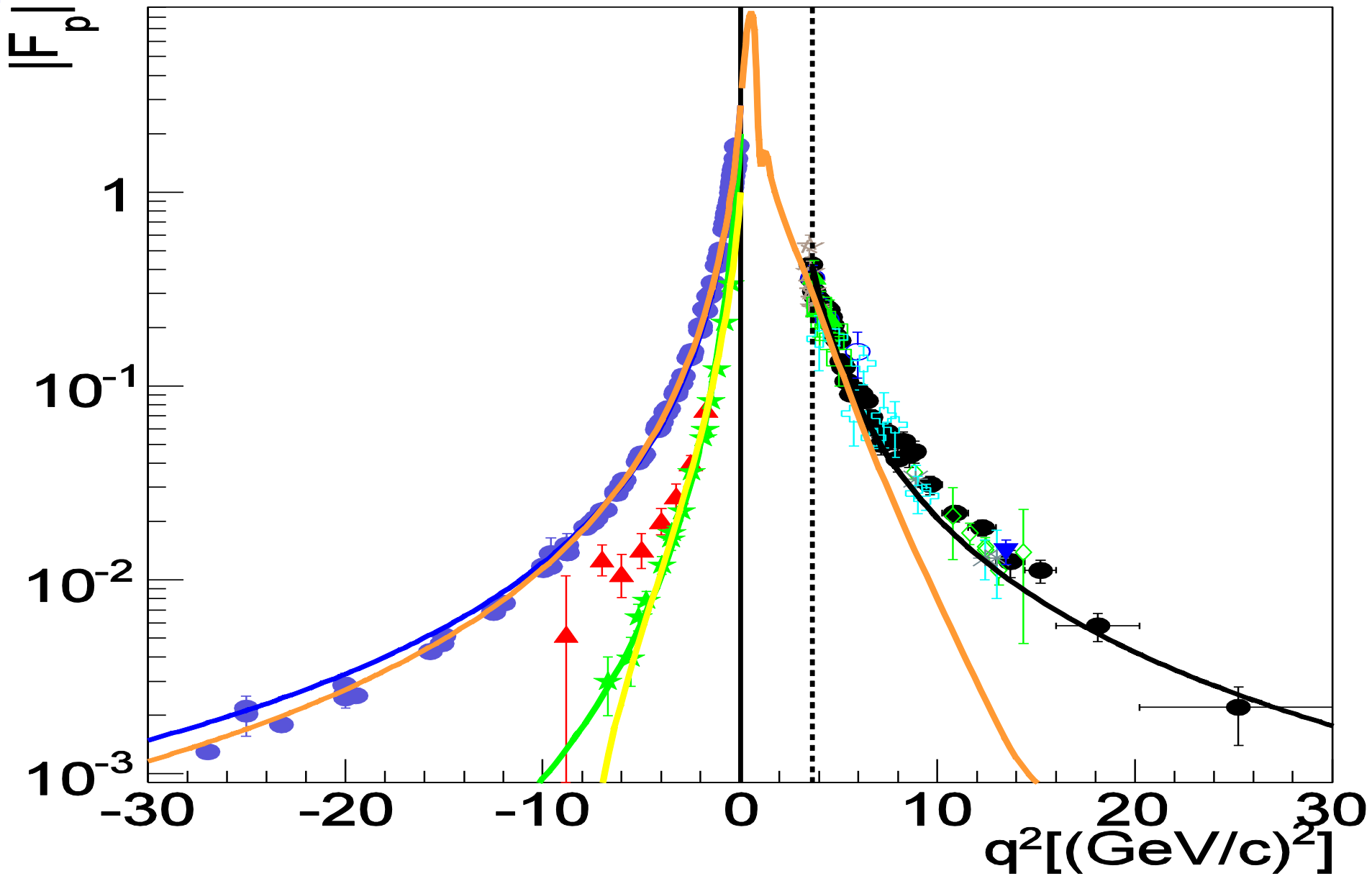
*Time-Like
FFs are complex*

$e+p \rightarrow e+p$

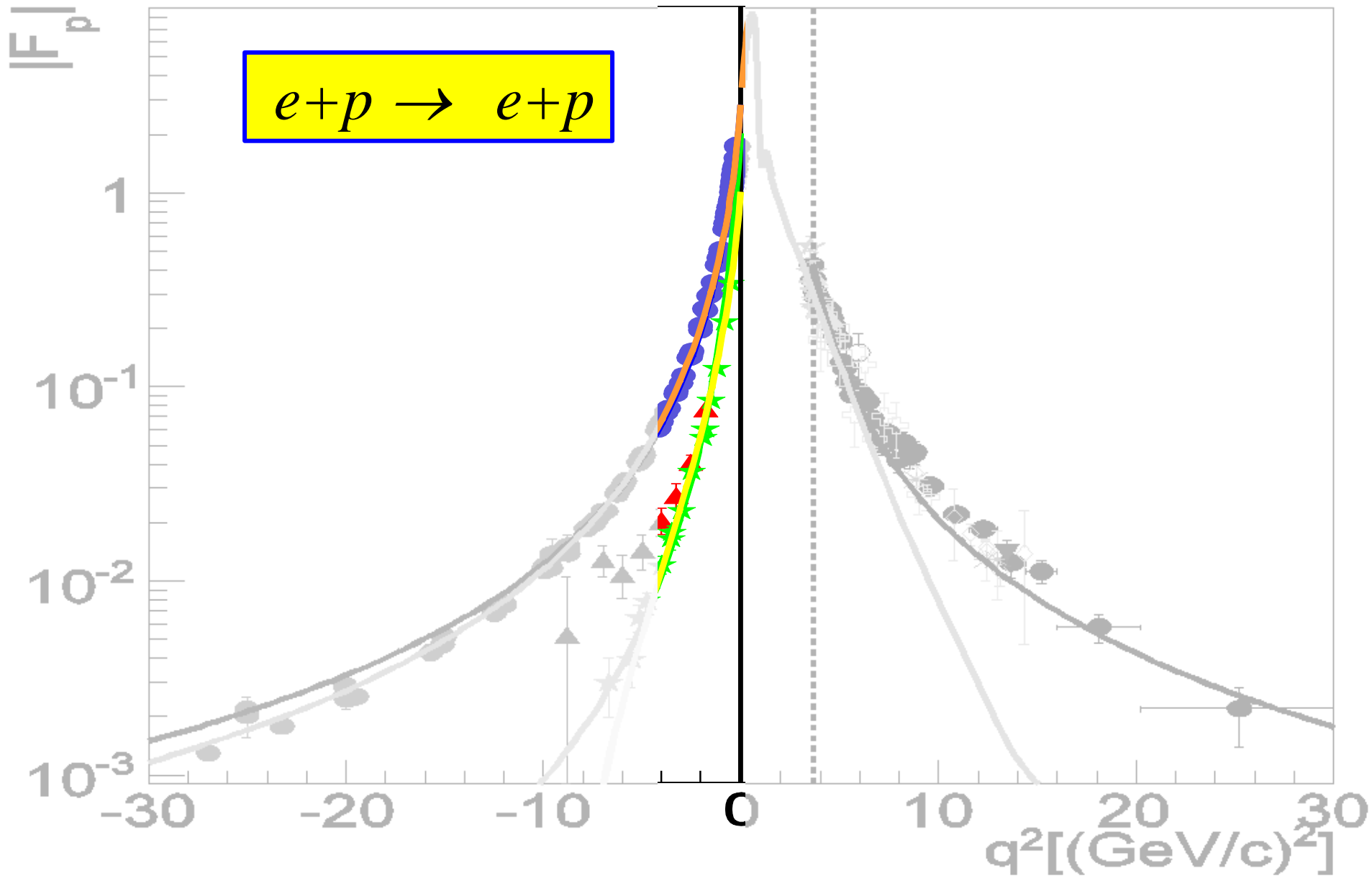
$q^2=4m_p^2$
 $GE=GM$

$p+\bar{p} \leftrightarrow e^+ + e^-$ q^2

Hadron Electromagnetic Form factors



The Space-Like region: low Q^2



The Proton Radius

$R_p=0.84184(67)$ fm (muonic atom)

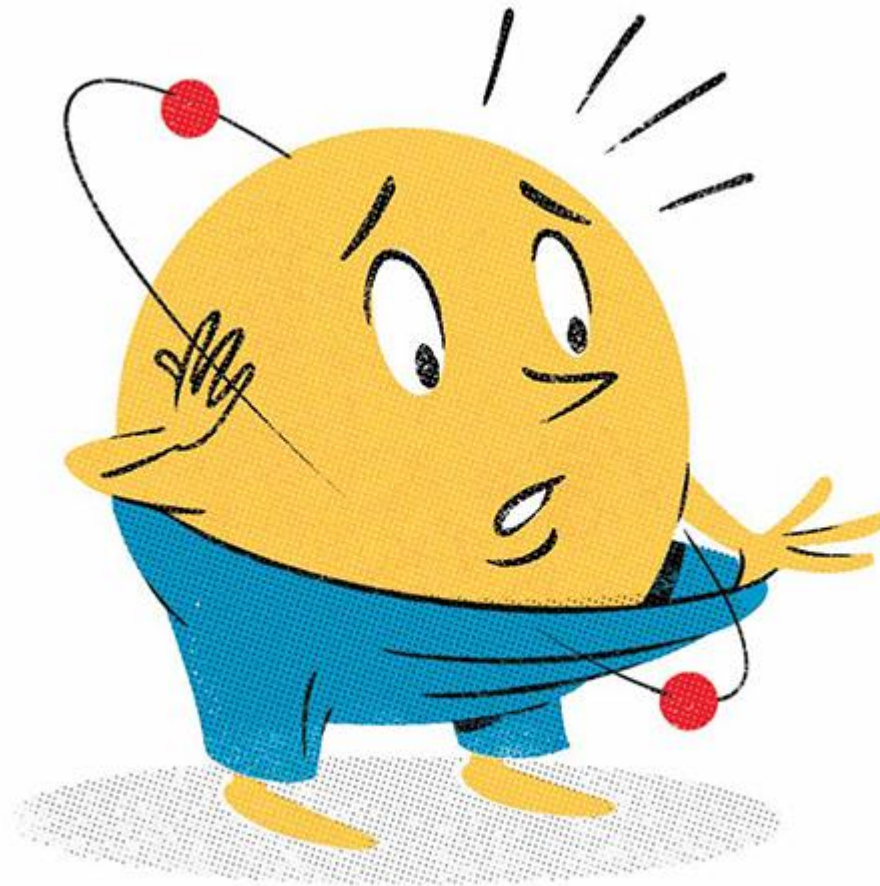
μp

dispersion

lattice QCD

$R_p=0.82-0.85$ fm
(DR PRC75 035202(2007))

$R_p=0.78-0.86$ fm
(lattice QCD PRD 79 094001(2009))



n

9) fm





High-Precision Determination of the Electric and Magnetic Form Factors of the Proton

J. C. Bernauer,^{1,*} P. Achenbach,¹ C. Ayerbe Gayoso,¹ R. Böhm,¹ D. Bosnar,² L. Debenjak,³ M. O. Distler,^{1,†} L. Doria,¹ A. Esser,¹ H. Fonvieille,⁴ J. M. Friedrich,⁵ J. Friedrich,¹ M. Gómez Rodríguez de la Paz,¹ M. Makek,² H. Merkel,¹ D. G. Middleton,¹ U. Müller,¹ L. Nungesser,¹ J. Pochodzalla,¹ M. Potokar,³ S. Sánchez Majos,¹ B. S. Schlimme,¹ S. Širca,^{6,3} Th. Walcher,¹ and M. Weinriefer¹

Mainz, A1 collaboration (1400 points)

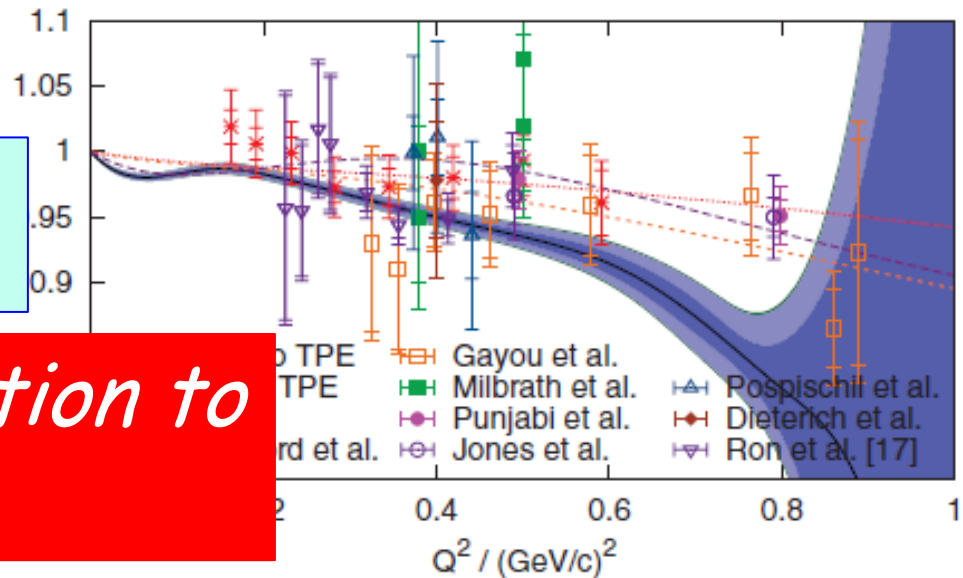
$Q^2 > 0.004 \text{ GeV}^2$

- Radiative corrections
- Two photon exchange
- Coulomb corrections

.....*comments*

$$\langle r_E^2 \rangle^{1/2} = 0.879(5)_{\text{stat}}(4)_{\text{syst}}(2)_{\text{model}}(4)_{\text{group}} \text{ fm},$$

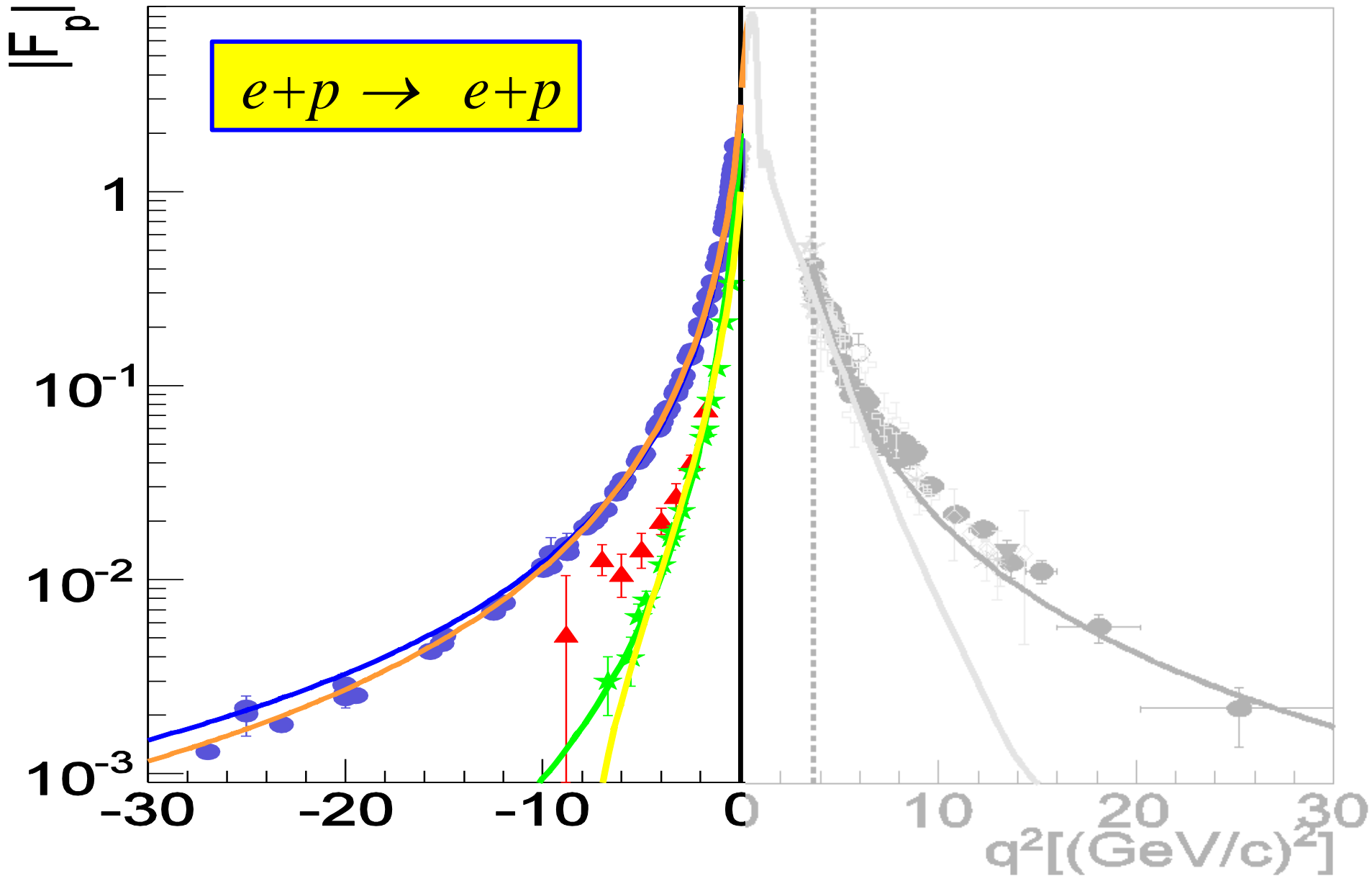
$$\langle r_M^2 \rangle^{1/2} = 0.777(13)_{\text{stat}}(9)_{\text{syst}}(5)_{\text{model}}(2)_{\text{group}} \text{ fm}.$$



- MUSE Experiment
- Jlab CLAS

What about extrapolation to $Q^2 \rightarrow 0$?

The Space-Like region

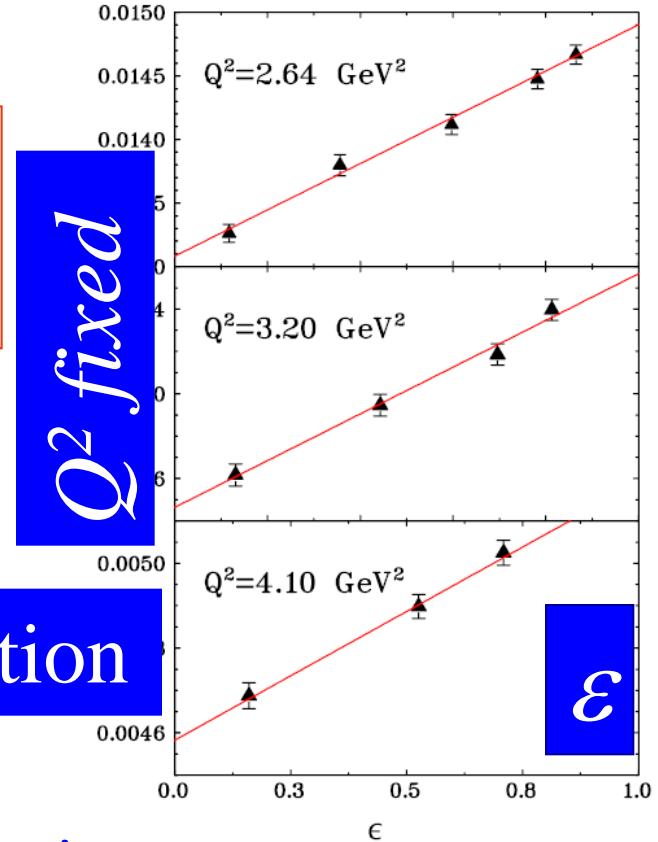


The Rosenbluth separation

$$\frac{d\sigma}{d\Omega} = \left(\frac{d\sigma}{d\Omega} \right)_{\text{Mott}} \frac{1}{(1+\tau)} \left(G_E^2(Q^2) + \frac{\tau}{\varepsilon} G_M^2(Q^2) \right)$$

$$\varepsilon = \left[1 + 2(1+\tau) \tan^2 \left(\frac{\theta_e}{2} \right) \right]^{-1}, \quad \tau = \frac{Q^2}{4M^2}$$

$$\sigma_R = \varepsilon G_E^2 + \tau G_M^2$$



Linearity of the reduced cross section

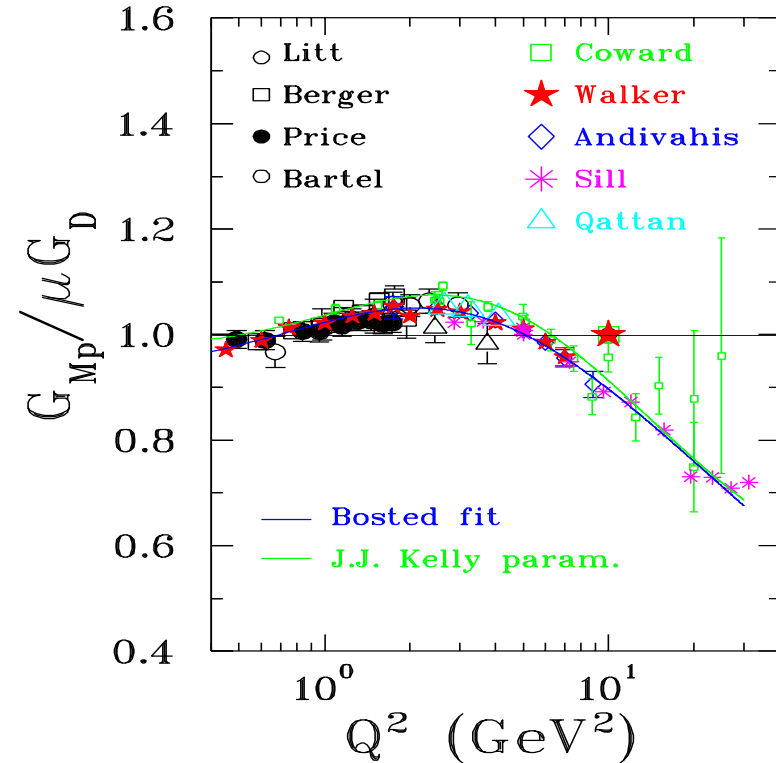
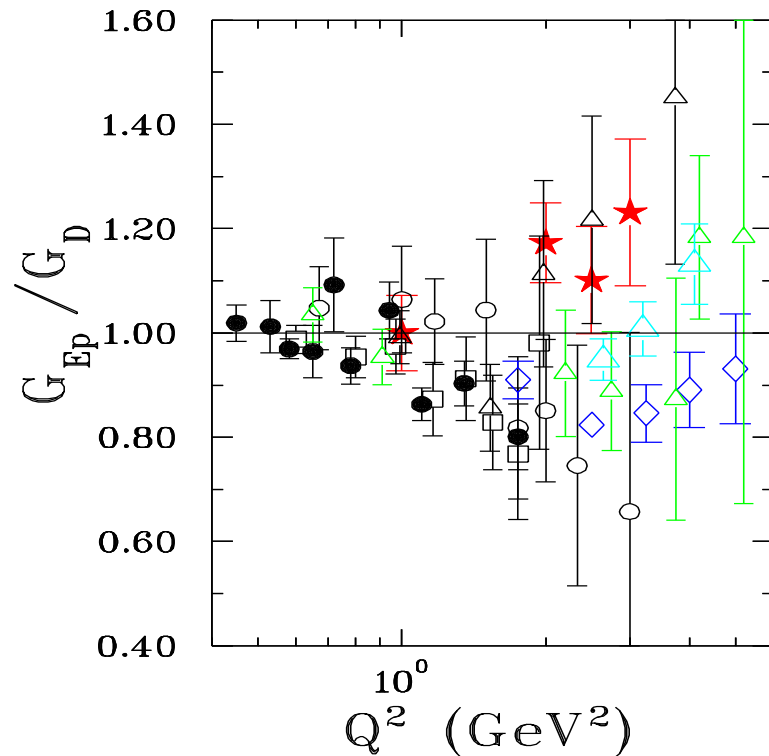
→ $\tan^2 \theta_e$ dependence

→ Holds for 1γ exchange only

PRL 94, 142301 (2005)

Proton Form Factors .. before

Dipole approximation: $G_D=(1+Q^2/0.71 \text{ GeV}^2)^{-2}$



Rosenbluth separation/ Polarization observables

V. Punjabi, M. Jones, C. Perdrisat et al, JLab-GEp collaboration

The polarization method (theory:1967)

SOVIET PHYSICS - DOKLADY

VOL. 13, NO. 6

DECEMBER, 1968

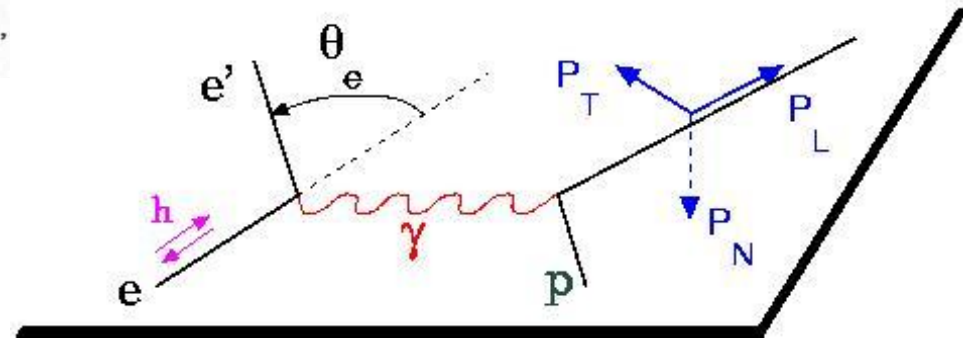
PHYSICS

POLARIZATION PHENOMENA IN ELECTRON SCATTERING BY PROTONS IN THE HIGH-ENERGY REGION

Academician A. I. Akhiezer* and M. P. Rekalo

Physicotechnical Institute, Academy of Sciences of the Ukrainian SSR
Translated from Doklady Akademii Nauk SSSR, Vol. 180, No. 5,
pp. 1081-1083, June, 1968
Original article submitted February 26,

$$s_2 \frac{d\sigma}{d\Omega_R} = 4p_2 \frac{(\mathbf{s} \cdot \mathbf{q})}{1 + \tau} \Gamma(\theta, \varepsilon_1) \left[\tau G_M (G_M + G_E) - \frac{1}{4\varepsilon_1} G_M (G_E - \tau G_M) \right],$$



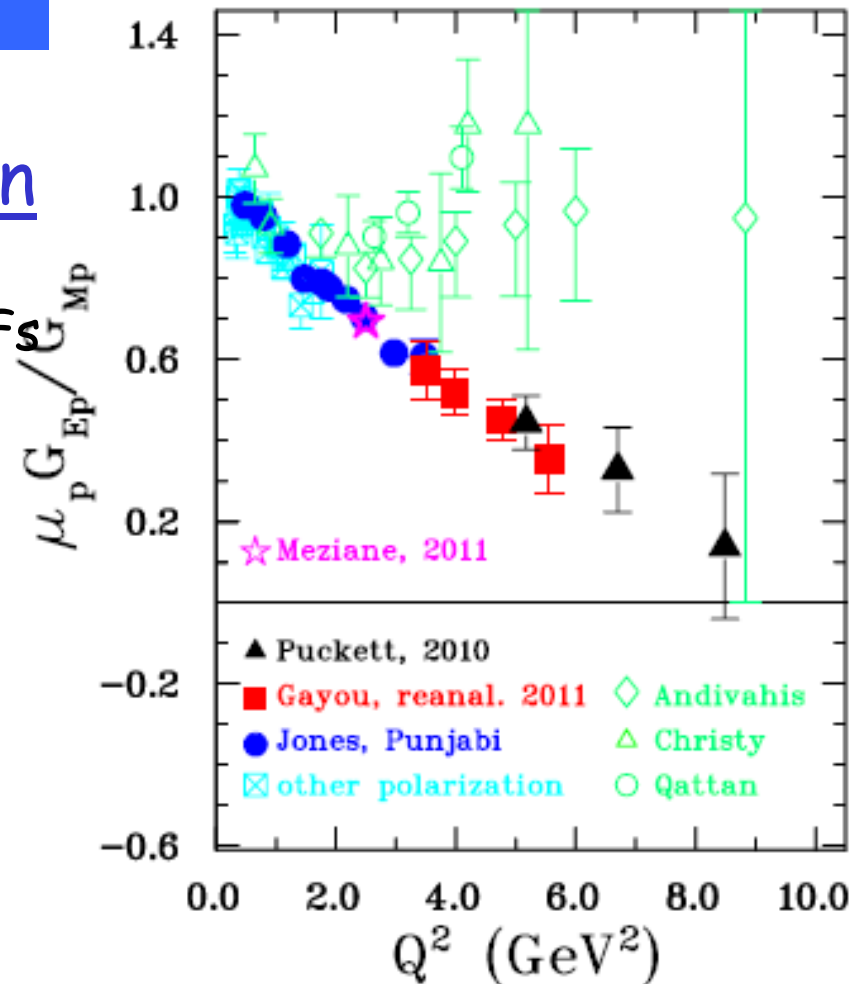
The polarization induces a term in the cross section proportional to $G_E G_M$
Polarized beam and target or
polarized beam and recoil proton polarization

Polarization experiments

A.I. Akhiezer and M.P. Rekalo, 1967

Jlab-GEp collaboration

- 1) "standard" dipole function for the nucleon magnetic FFs G_{Mp} and G_{Mn}
- 2) linear deviation from the dipole function for the electric proton FF G_{Ep}
- 3) QCD scaling not reached
- 3) Zero crossing of G_{Ep} ?
- 4) contradiction between polarized and unpolarized measurements



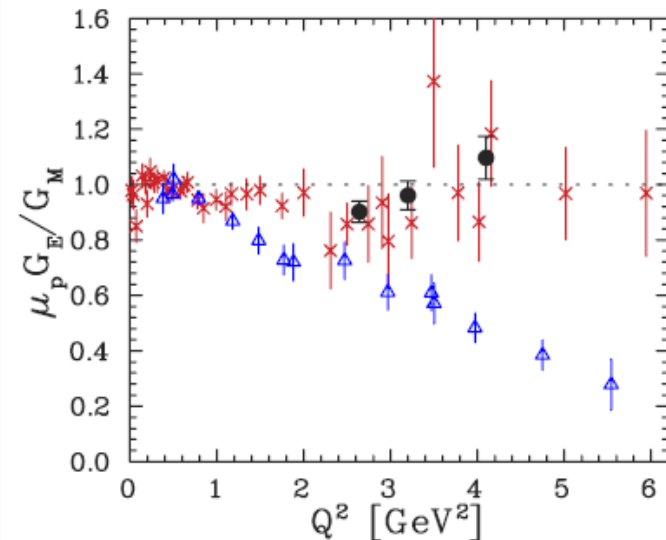
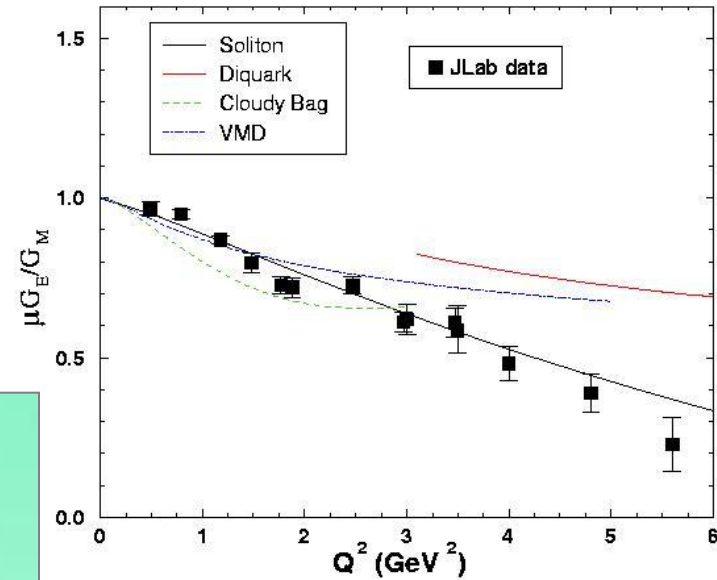
A.J.R. Puckett et al, PRL (2010), PRC (2012)

ISSUES

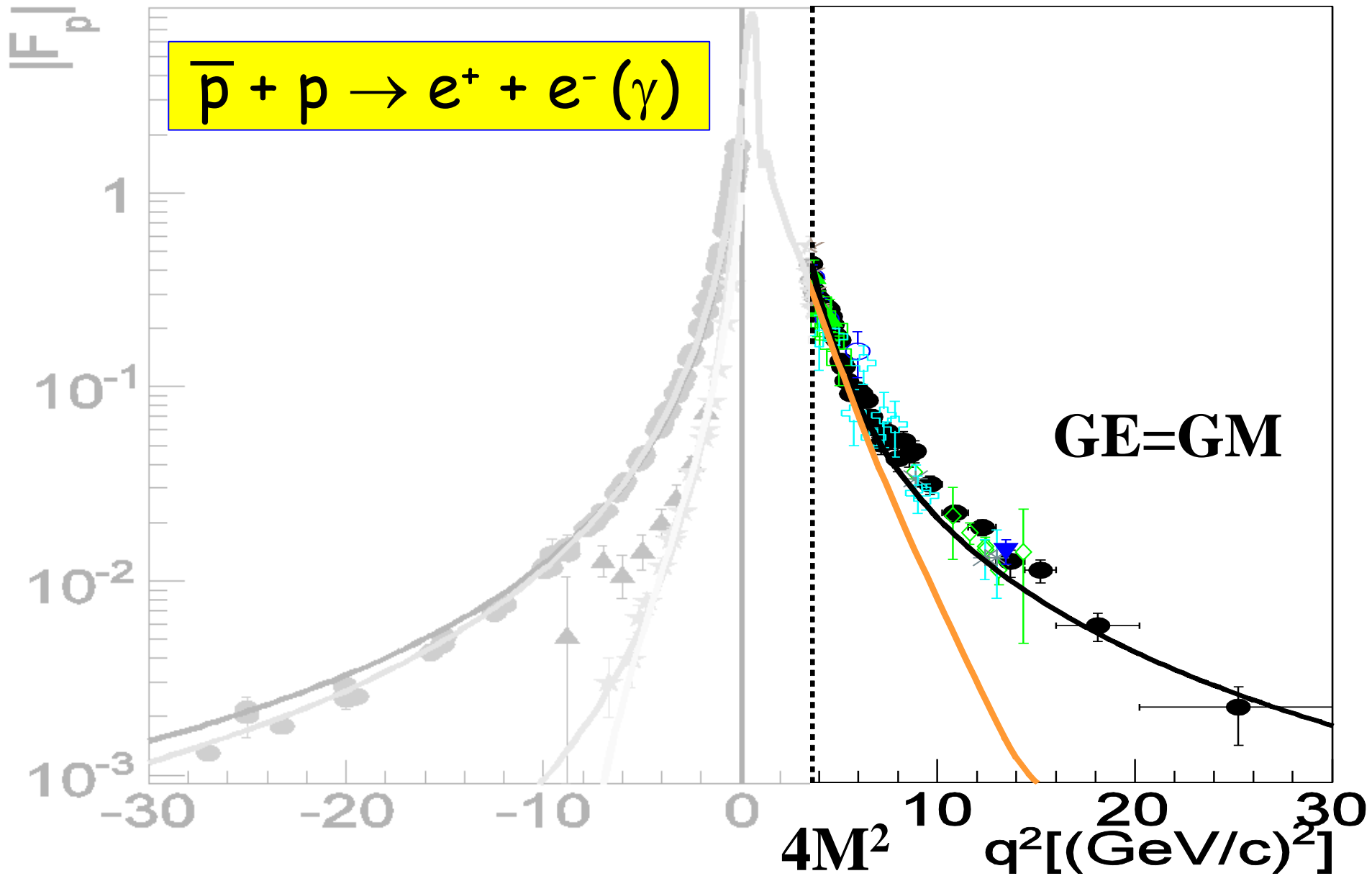
- Some models (IJL 73, Diquark, soliton..) predicted such behavior before the data appeared

BUT

- Simultaneous description of the four nucleon form factors...
- ...in the space-like and in the time-like regions
- Consequences for the light ions description
- When pQCD starts to apply?
- Source of the discrepancy



The Time-Like region



Time-like observables: $|G_E|^2$ and $|G_M|^2$.

-The cross section for $\bar{p} + p \rightarrow e^+ + e^-$ (1 γ -exchange):

$$\frac{d\sigma}{d(\cos \theta)} = \frac{\pi\alpha^2}{8m^2\sqrt{\tau-1}} [\tau|G_M|^2(1 + \cos^2 \theta) + |G_E|^2 \sin^2 \theta]$$

θ : angle between e^- and \bar{p} in cms.

A. Zichichi, S. M. Berman, N. Cabibbo, R. Gatto, Il Nuovo Cimento XXIV, 170 (1962)

B. Bilenkii, C. Giunti, V. Wataghin, Z. Phys. C 59, 475 (1993).

G. Gakh, E.T-G., Nucl. Phys. A761,120 (2005).

As in SL region:

- Dependence on q^2 contained in FFs
- Even dependence on $\cos^2\theta$ (1 γ exchange)
- No dependence on sign of FFs
- Enhancement of magnetic term

but TL form factors are complex!

The Experimental facilities

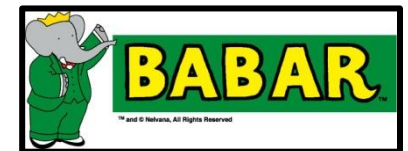
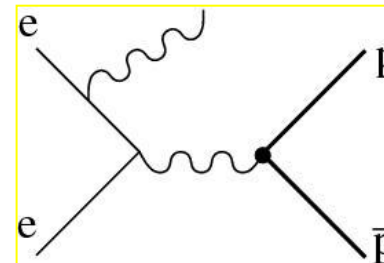
- Antiproton-proton colliders:
 - LEAR, FERMILAB, PANDA
- Electron -positron colliders
 - FENICE, VEPP, BABAR, BES
- Initial State Radiation
 - BABAR, BES



VEPP-Novosibirsk



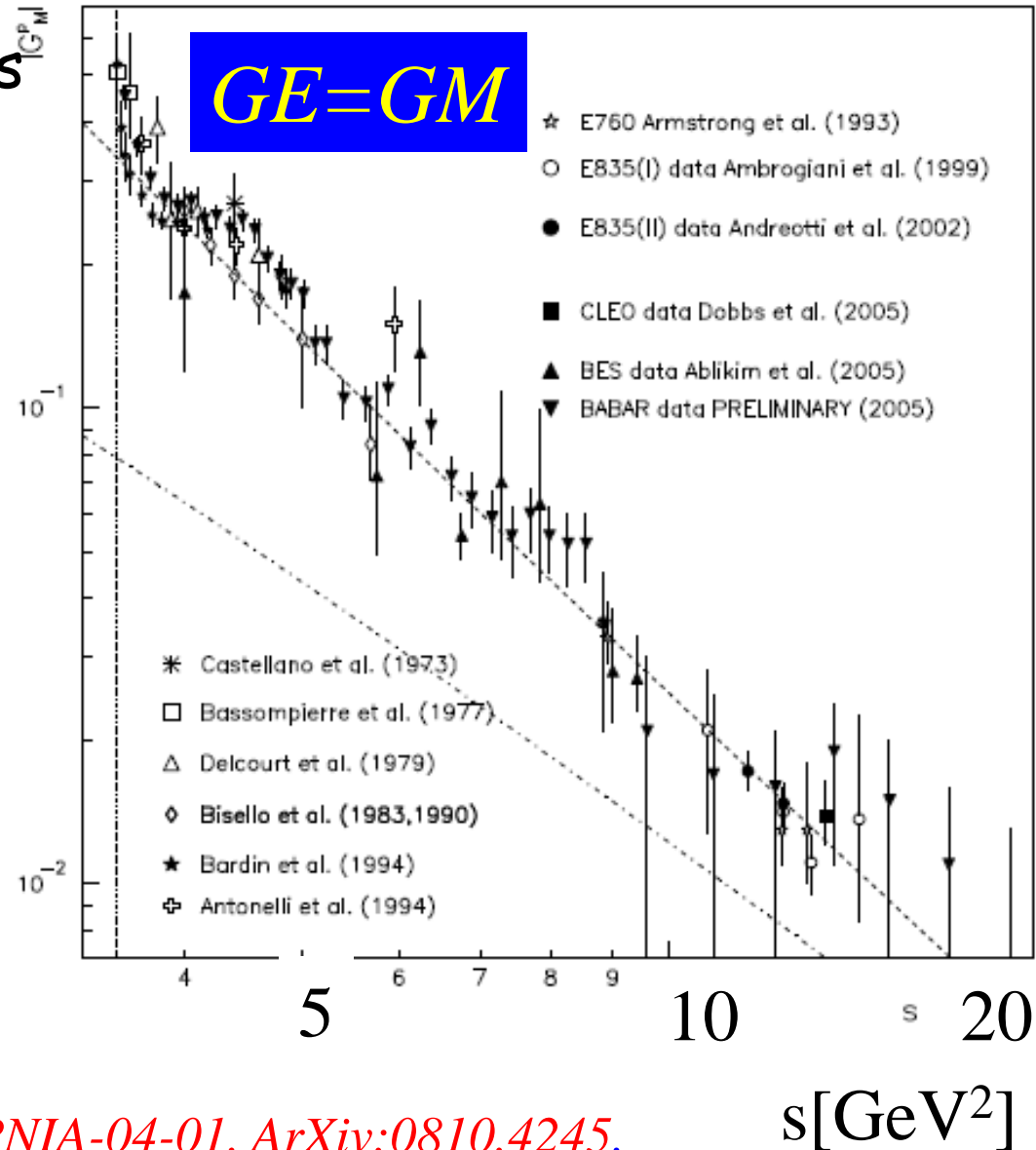
BES



The Time-like region

The Experimental Status^[CPM]

- No individual determination of GE and GM
- TL proton FFs twice larger than in SL at the same Q^2
- Steep behaviour at threshold
- Babar:
 - Structures?
 - Resonances?
- -> Panda contribution



MP. Rekaló, E.T-G., preprint DAPNIA-04-01, ArXiv:0810.4245.

Antiproton facilities

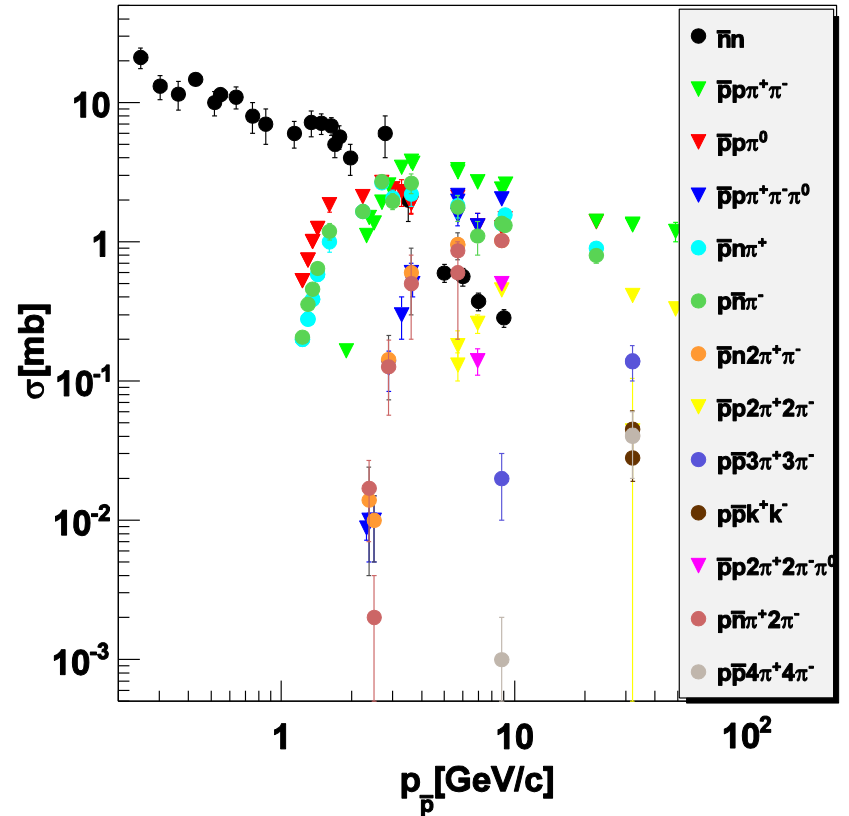
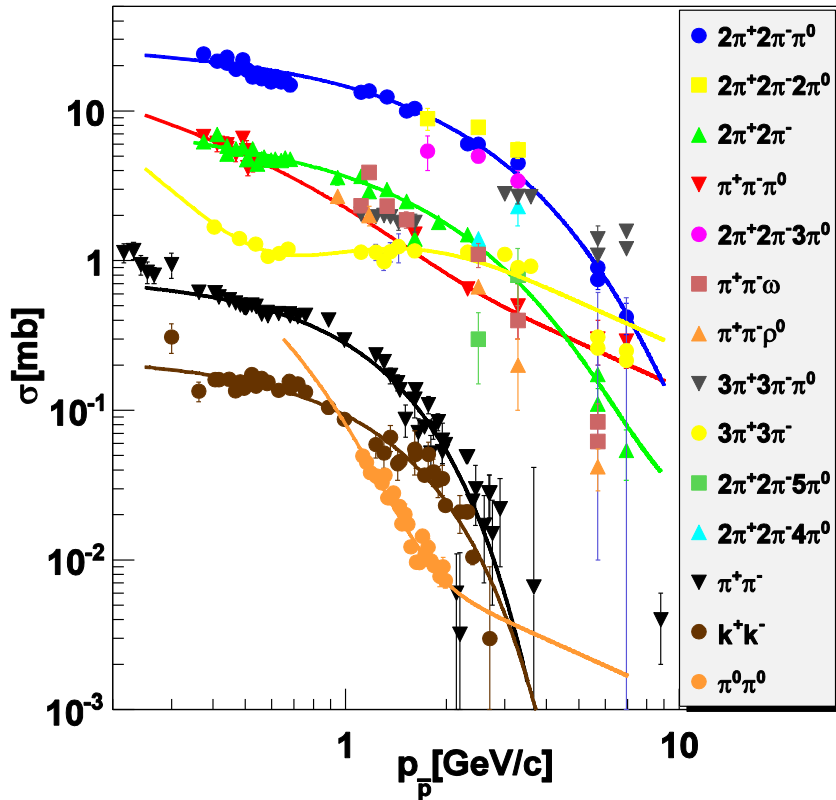
Experiment	Years	Intensity \bar{p}/s	Momentum range [GeV/c]	$\Delta p/p$
CERN -LEAR	1983-1996	$2 \cdot 10^6$	0.06-1.94	10^{-3}
FermiLab 45% polarized \bar{p}	1985-2011	$2 \cdot 10^6$ 10^4	<8.9 (Low energy beams)	10^{-4}
PANDA		$2 \cdot 10^7$	1.5-15	10^{-5}

Panda will have:

- Better luminosity*
- Better beam momentum resolution*
- Better detector (coverage, PID, magnetic field..)*

About cross sections...

A.Dbeyssi and E.T.-G, *Prob Atomic Sci. Technol.* 2012N1, 79 (2012).



Most probable in the mb region :

- *Five pion production*
- *Charge exchange*

The following slides concerning ongoing PANDA simulations :

- are extracted/derived from:

A. Dbeyssi, PhD thesis
Université Paris XI, 27-IX-2013

- they are considered
'PANDA unofficial results'

Past its prime PANDA simulations are published in:

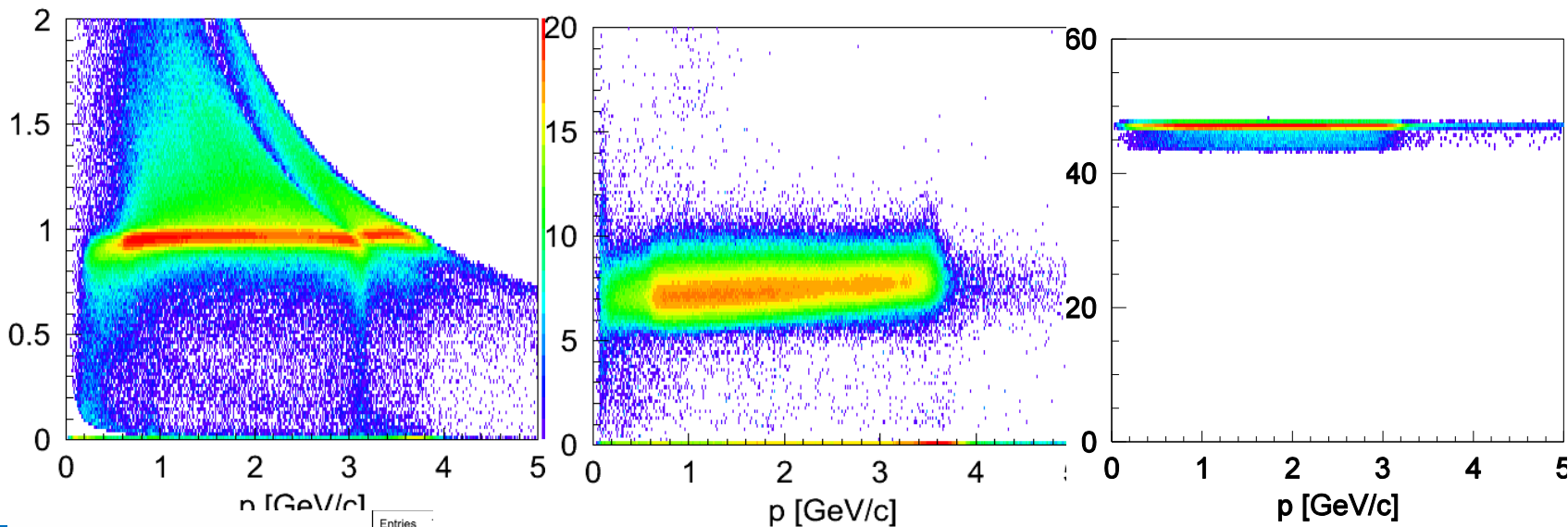
- *M.C. Mora Espi, PhD Mainz 2013*
- *M. Sudol et al., EPJA A44(2010)373*

EMCE_{raw}/p vs. p

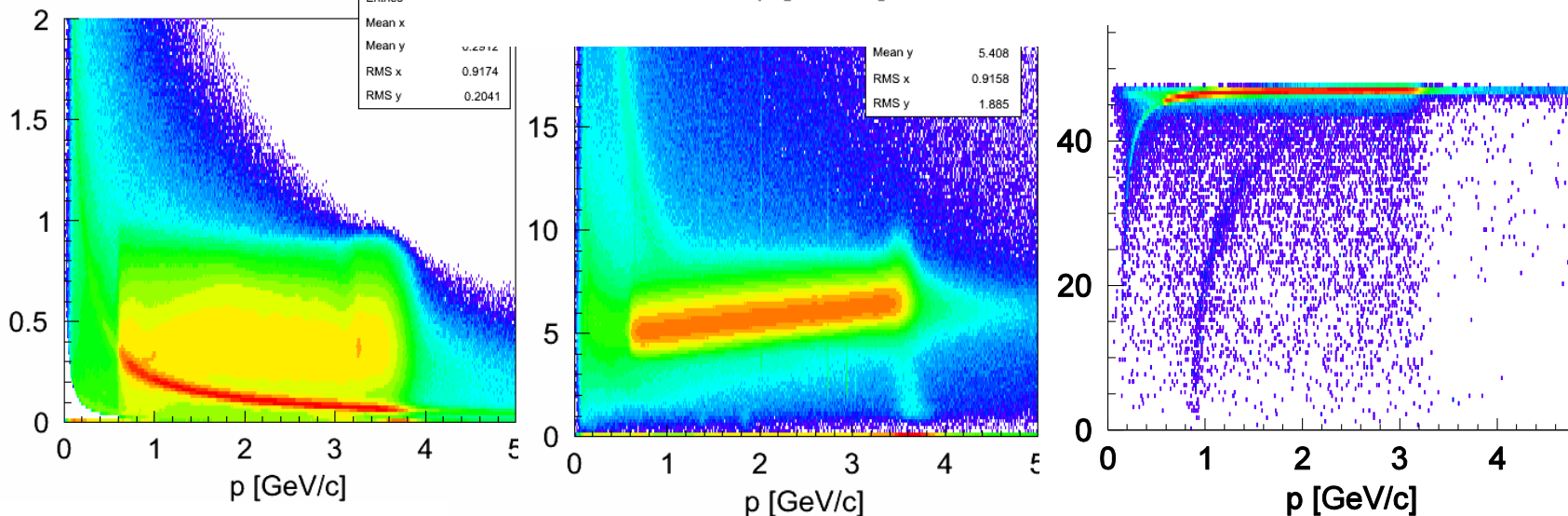
STT dE/dx vs p

Barrel DIRC Θ Cherenkov

SIGNAL



BACKGROUND



From PHSP to physical angular distributions

A.DBEYSSI

PhD-2013

The differential cross section in the CM for $\bar{p}p \rightarrow e^+e^-$ is:

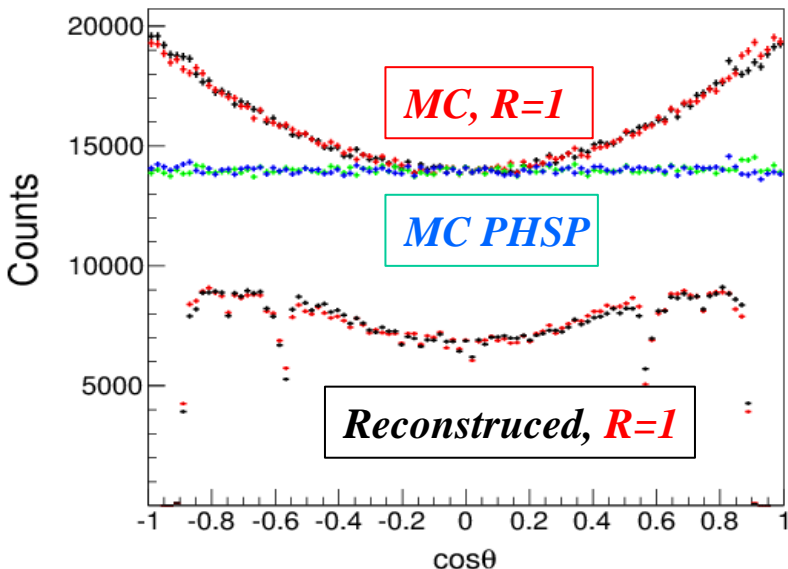
$$\frac{d\sigma}{d\cos\theta} = \sigma_0(1 + \mathcal{A} \cos^2\theta)$$

$$\left\{ \begin{array}{l} \sigma_0 = \frac{d\sigma}{d\cos\theta} \left(\theta = \frac{\pi}{2} \right) \\ \mathcal{A} = \frac{\tau - R^2}{\tau + R^2}, R = |G_E|/|G_M|, \tau = \frac{s}{4M^2} \end{array} \right.$$

$$\mathcal{A} = \frac{\tau - R^2}{\tau + R^2}, R = |G_E|/|G_M|, \tau = \frac{s}{4M^2}$$

A. Zichichi et al., *Nuovo Cim.* 24 (1962) 170

E. Tomasi-Gustafsson and M.P. Rekalo, *Phys.Lett.* B504 (2001) 291-295



Monte Carlo events, PHSP

× Weight: $1 + \mathcal{A} \cos^2\theta$

Physical Monte Carlo events

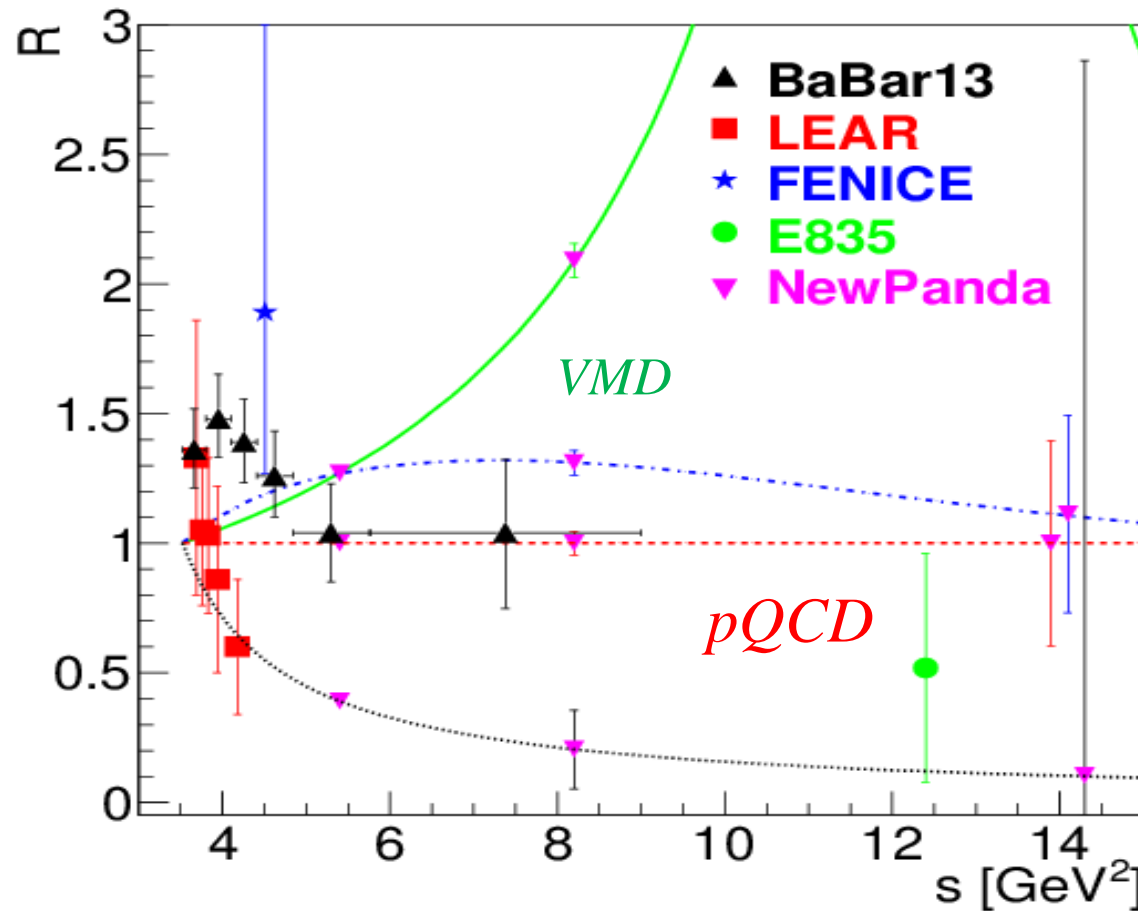
× Efficiency $\epsilon(c)$

Physical reconstructed events

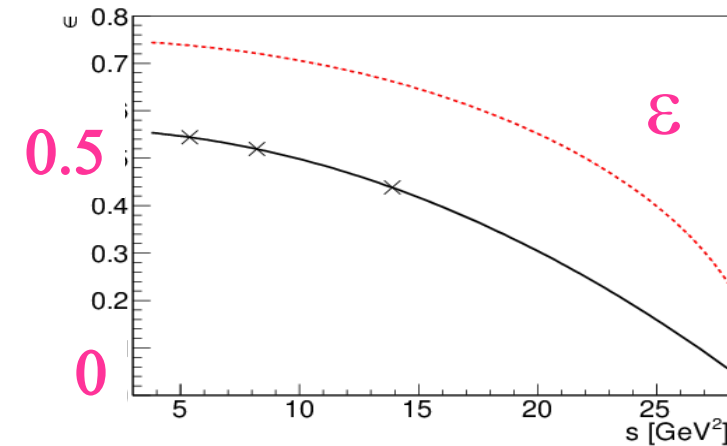
Individual determination of $|G_E|$ and $|G_M|$

A.Dbeyssi

PhD 2013

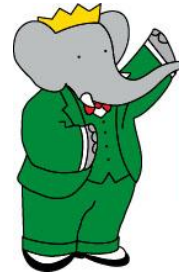
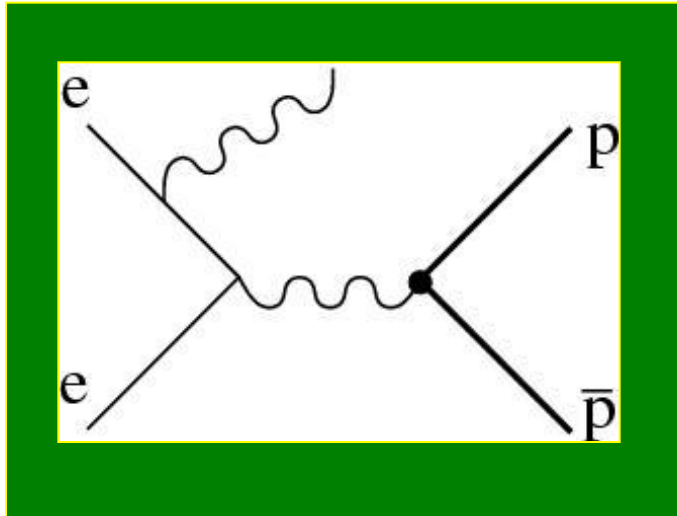


$\epsilon \leq 50\%$



F. Iachello et al., Phys. Rev. C 69 (2004) 055204 *E. A. Kuraev et al., Phys. Lett. B 712, (2012)*
E. L. Lomon, Phys. Rev. C 66 (2002) 045501 *V. A. Matveev, S. J. Brodsky, D. V. Shirkov....*

Radiative return (ISR)



BABARTM

TM and © Neivana, All Rights Reserved

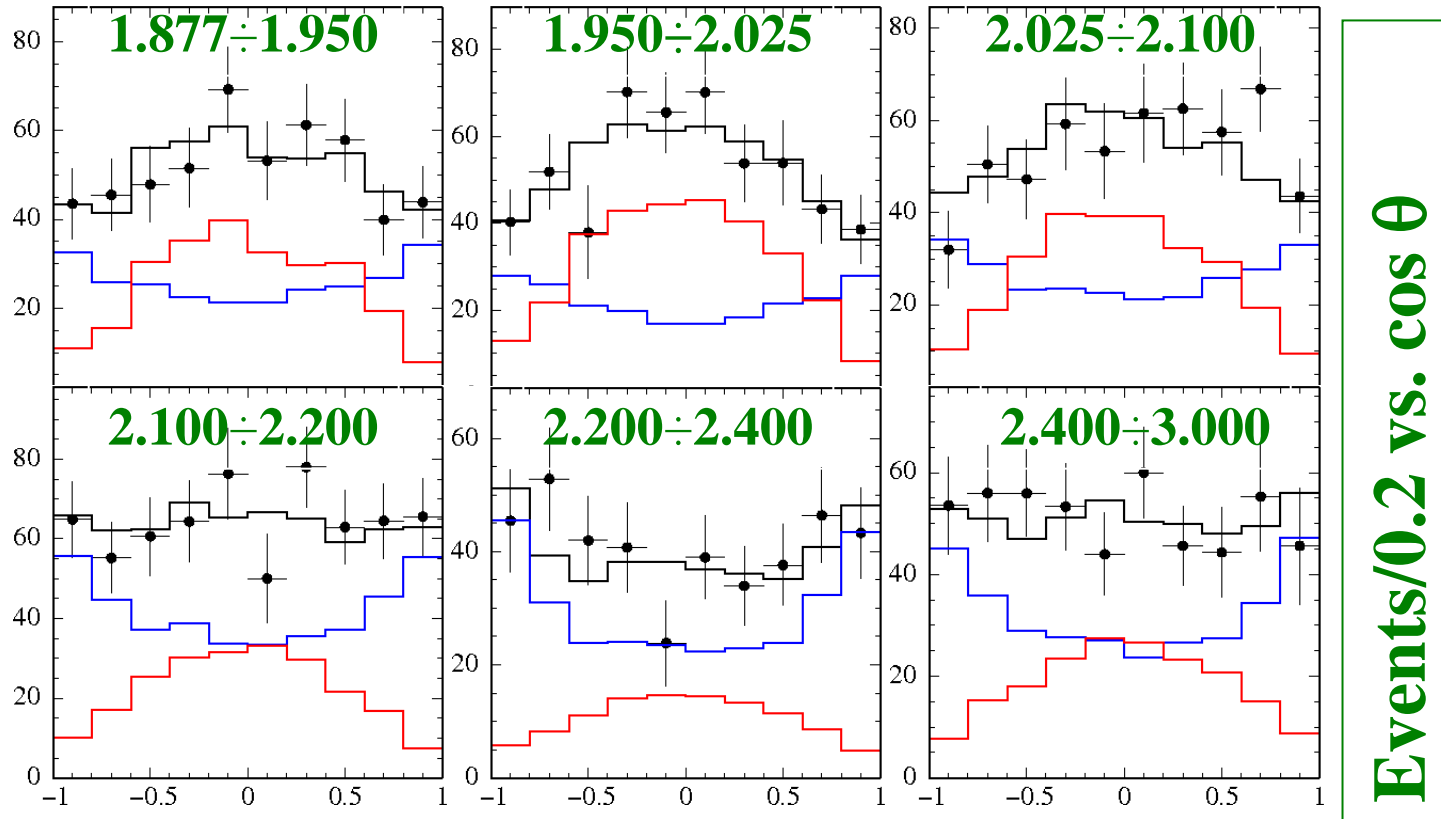


$$\frac{d\sigma(e^+e^- \rightarrow p\bar{p}\gamma)}{dm d\cos\theta} = \frac{2m}{s} W(s, x, \theta) \sigma(e^+e^- \rightarrow p\bar{p})(m), \quad x = \frac{2E_\gamma}{\sqrt{s}} = 1 - \frac{m^2}{s},$$

$$W(s, x, \theta) = \frac{\alpha}{\pi x} \left(\frac{2 - 2x + x^2}{\sin^2\theta} - \frac{x^2}{2} \right), \quad \theta \gg \frac{m_e}{\sqrt{s}}.$$

B. Aubert (BABAR Collaboration) Phys Rev. D73, 012005 (2006)

Angular distribution



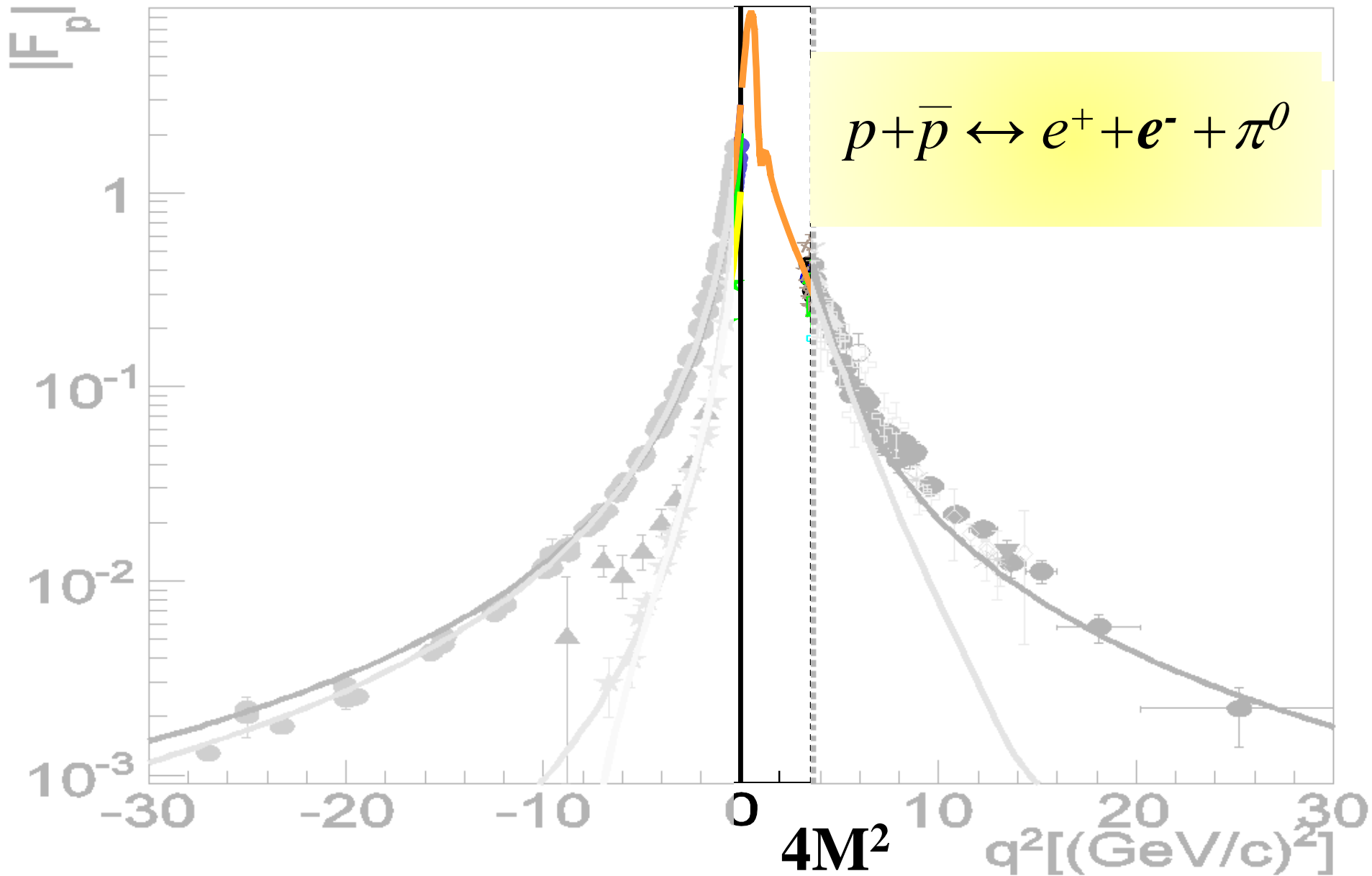
$$\frac{dN}{d \cos \theta_p} = A \left[H_M(\cos \theta, M_{pp}) - \left| \frac{G_E}{G_M} \right|^2 H_E(\cos \theta, M_{pp}) \right]$$



BABAR

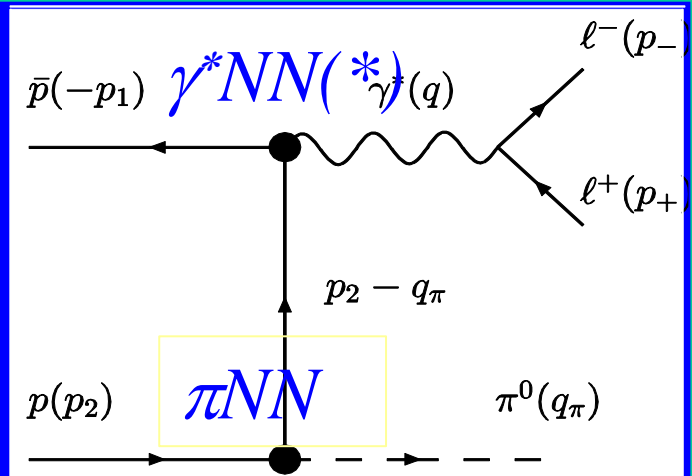
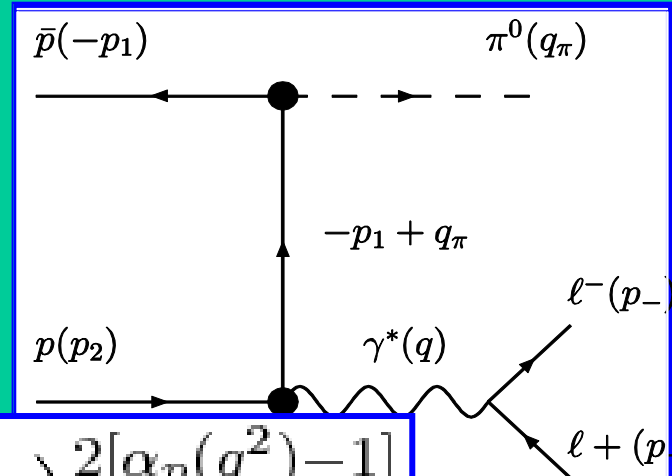
™ and © Helvex, All Rights Reserved

The "unphysical region"



The reaction $p + \bar{p} \rightarrow e^+ + e^- + \pi^0$

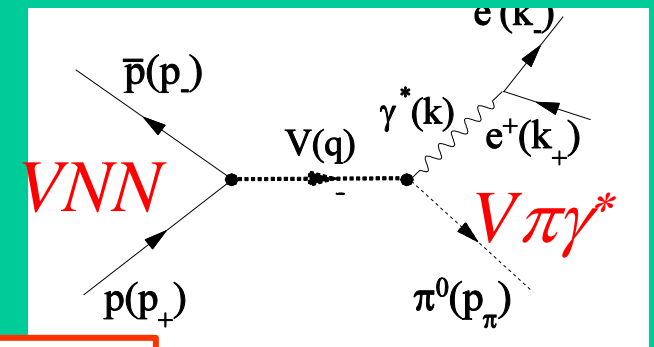
*E.A. Kuraev,
PRC, 2008
JETP, 2012
G.I.Gakh et al,
PRC, 2011*



$$d\sigma_s \propto \frac{1}{q^2} \left(\frac{s}{M^2} \right)^{2[\alpha_p(q^2) - 1]}$$

M. P. Rekalo, 1967

$V = \rho, \omega, \phi, J/\Psi, \dots$



$$d\sigma_a \propto 1/s$$

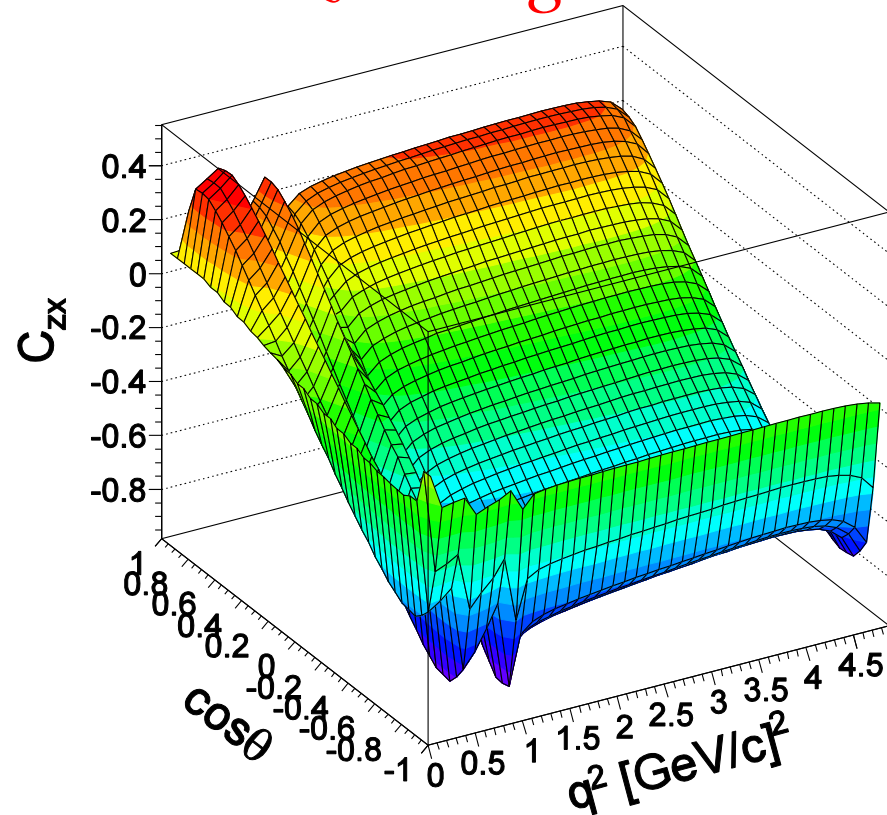
Forward/backward
Regge trajectory
of the proton

$$\frac{d\sigma_s}{d\sigma_a} \propto \left(\frac{s}{M^2} \right)^{2[\alpha_p(q^2) - 1]} \ll 1$$

Large angle: VM exchange

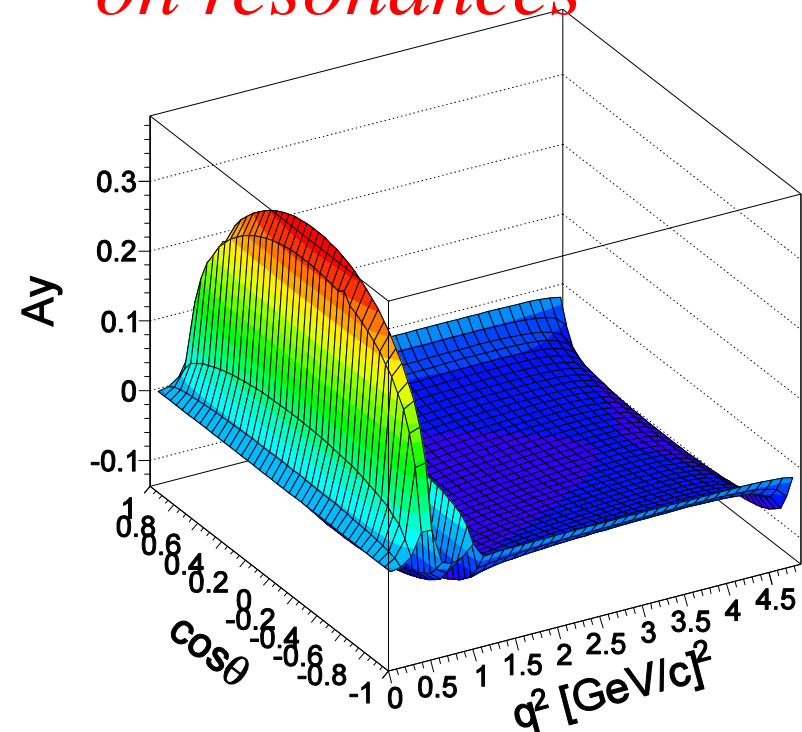
Results with IJL FFs

C_{zx}: large at small angles



$2E=5.4 \text{ GeV}$

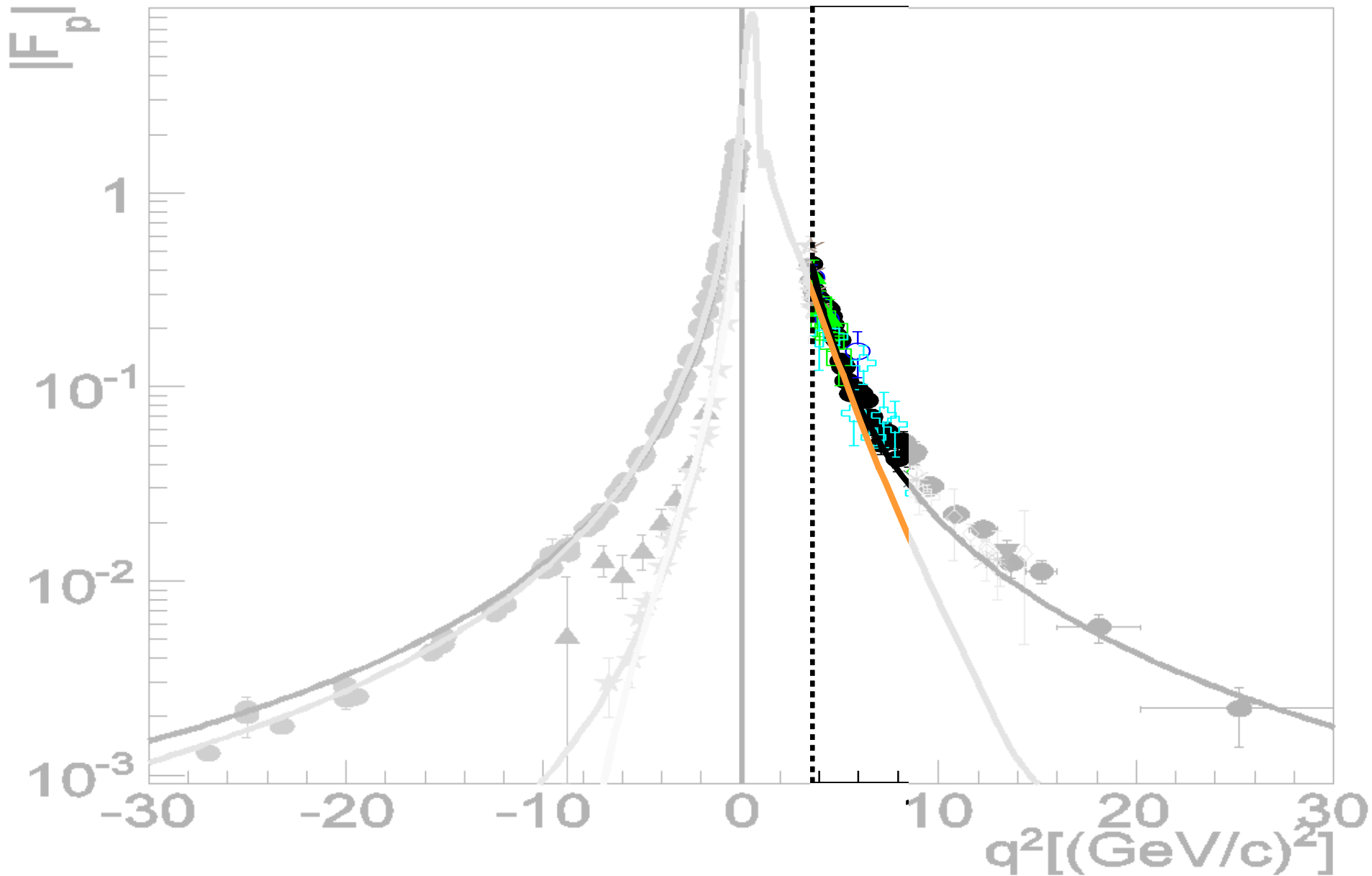
A_y: small except on resonances



*Polarization phenomena help
to distinguish the reaction mechanism!*

G.I. Gakh, J. Boucher, E.T-G., Phys.Rev. C83 (2011)

The Time-like region: the threshold



Point-like form factors?

Sommerfeld Enhancement and Resummation Factors

S. Pacetti

Coulomb Factor \mathcal{C} for S-wave only:

● Partial wave FF: $G_S = \frac{2G_M \sqrt{q^2/4M^2} + G_E}{3}$ $G_D = \frac{G_M \sqrt{q^2/4M^2} - G_E}{3}$

● Cross section: $\sigma(q^2) = 2\pi\alpha^2\beta \frac{4M^2}{(q^2)^2} \left[\mathcal{C} |G_S(q^2)|^2 + 2|G_D(q^2)|^2 \right]$

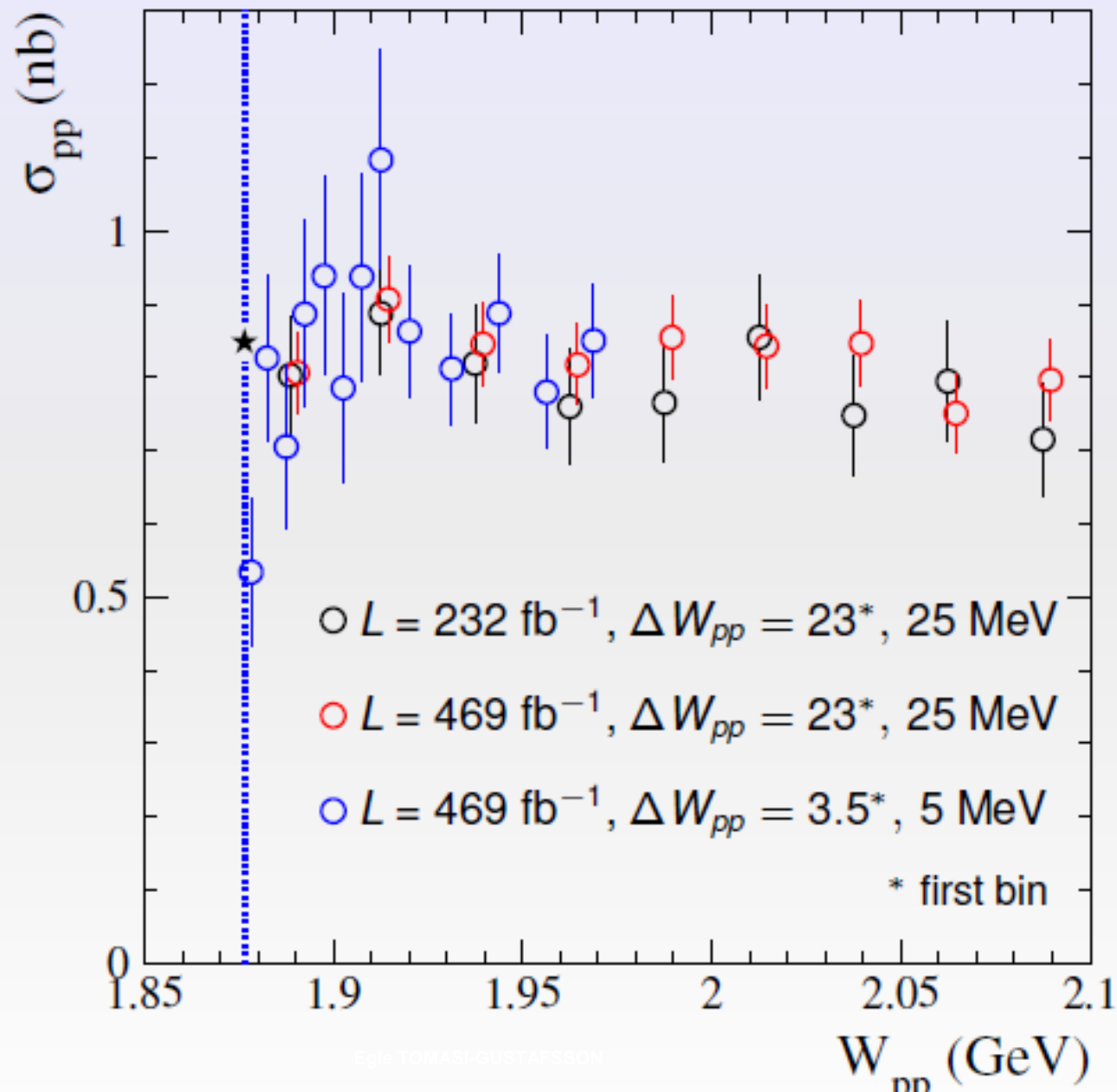
$$\mathcal{C} = \mathcal{E} \times \mathcal{R}$$

● Enhancement factor: $\mathcal{E} = \pi\alpha/\beta$

● Step at threshold: $\sigma_{p\bar{p}}(4M_p^2) = \frac{\pi^2\alpha^3}{2M^2} \frac{\beta}{\beta} |G_S^p(4M_p^2)|^2 = 0.85 |G_S^p(4M_p^2)|^2 \text{ nb}$

● Resummation factor: $\mathcal{R} = 1/[1 - \exp(-\pi\alpha/\beta)]$

● Few MeV above threshold: $\mathcal{C} \simeq 1 \Rightarrow \sigma_{p\bar{p}}(q^2) \propto \beta |G_S^p(q^2)|^2$



The nucleon



*3 valence quarks and
a neutral sea of $\bar{q}q$ pairs*

*antisymmetric state of
colored quarks*

$$|p\rangle \sim \epsilon_{ijk} |u^i u^j d^k\rangle$$
$$|n\rangle \sim \epsilon_{ijk} |u^i d^j d^k\rangle$$

Main assumption

Does not hold in the spatial center of the nucleon: the center of the nucleon *is electrically neutral*, due to strong gluonic field

E.A. Kuraev, E. T-G, A. Dbeyssi, Phys.Lett. B712 (2012) 240

Model: generalized form factors

E.A. Kuraev, E. T-G, A. Dbeyssi, Phys. Lett. B712 (2012) 240

Definition:

$$F(q^2) = \int_{\mathcal{D}} d^4x e^{iq_\mu x^\mu} \rho(x), \quad q_\mu x^\mu = q_0 t - \vec{q} \cdot \vec{x}$$

$\rho(x) = \rho(\vec{x}, t)$ *space-time distribution of the electric charge in the space-time volume \mathcal{D} .*

In SL-Breit frame (zero energy transfer):

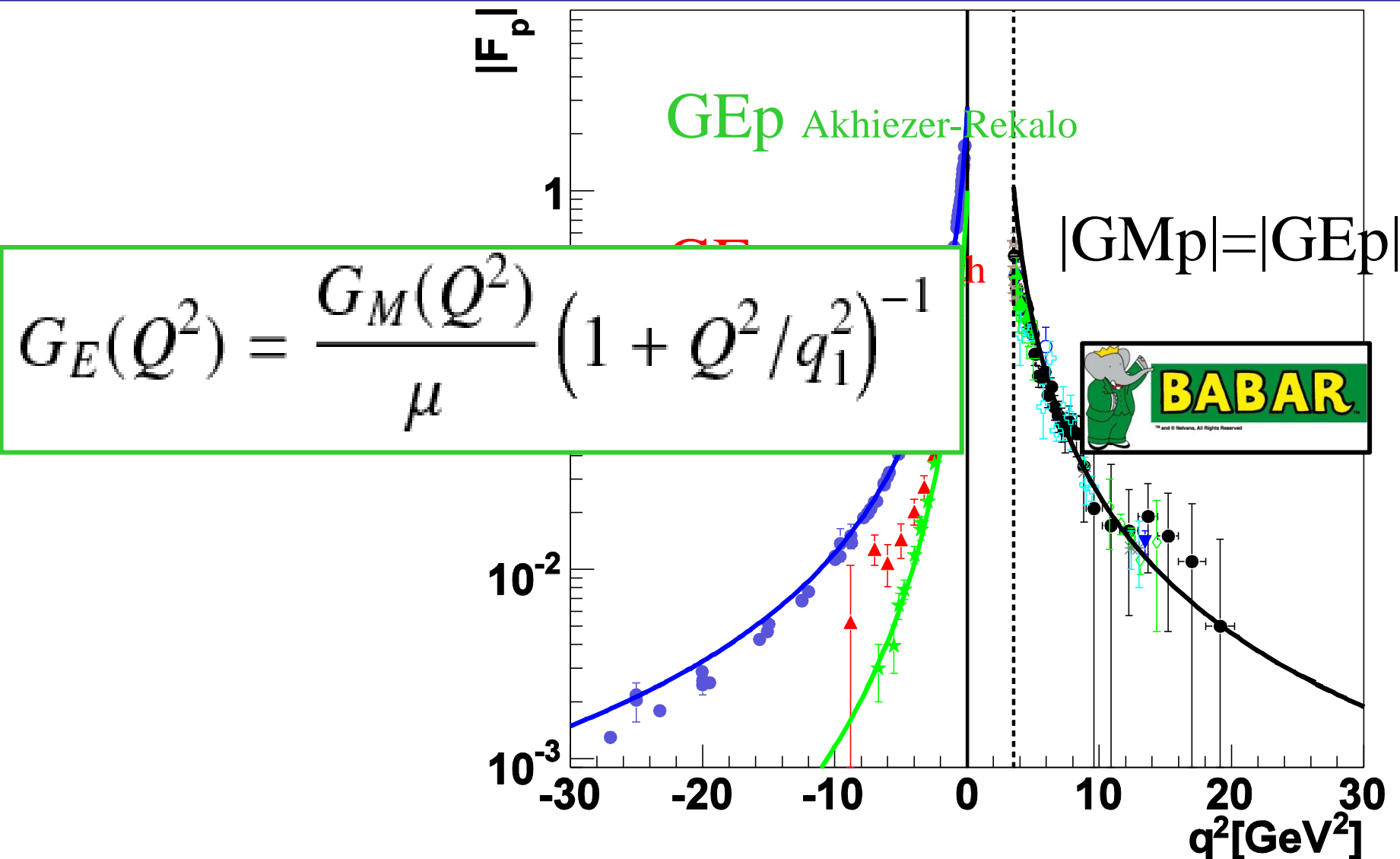
$$F(q^2) = \delta(q_0) F(Q^2), \quad Q^2 = -(q_0^2 - \vec{q}^2) > 0.$$

In TL-(CMS):

$$F(q^2) = \int_{\mathcal{D}} dt e^{i\sqrt{q^2}t} \int d^3\vec{r} \rho(\vec{r}, t) = \int_{\mathcal{D}} dt e^{i\sqrt{q^2}t} Q(t),$$

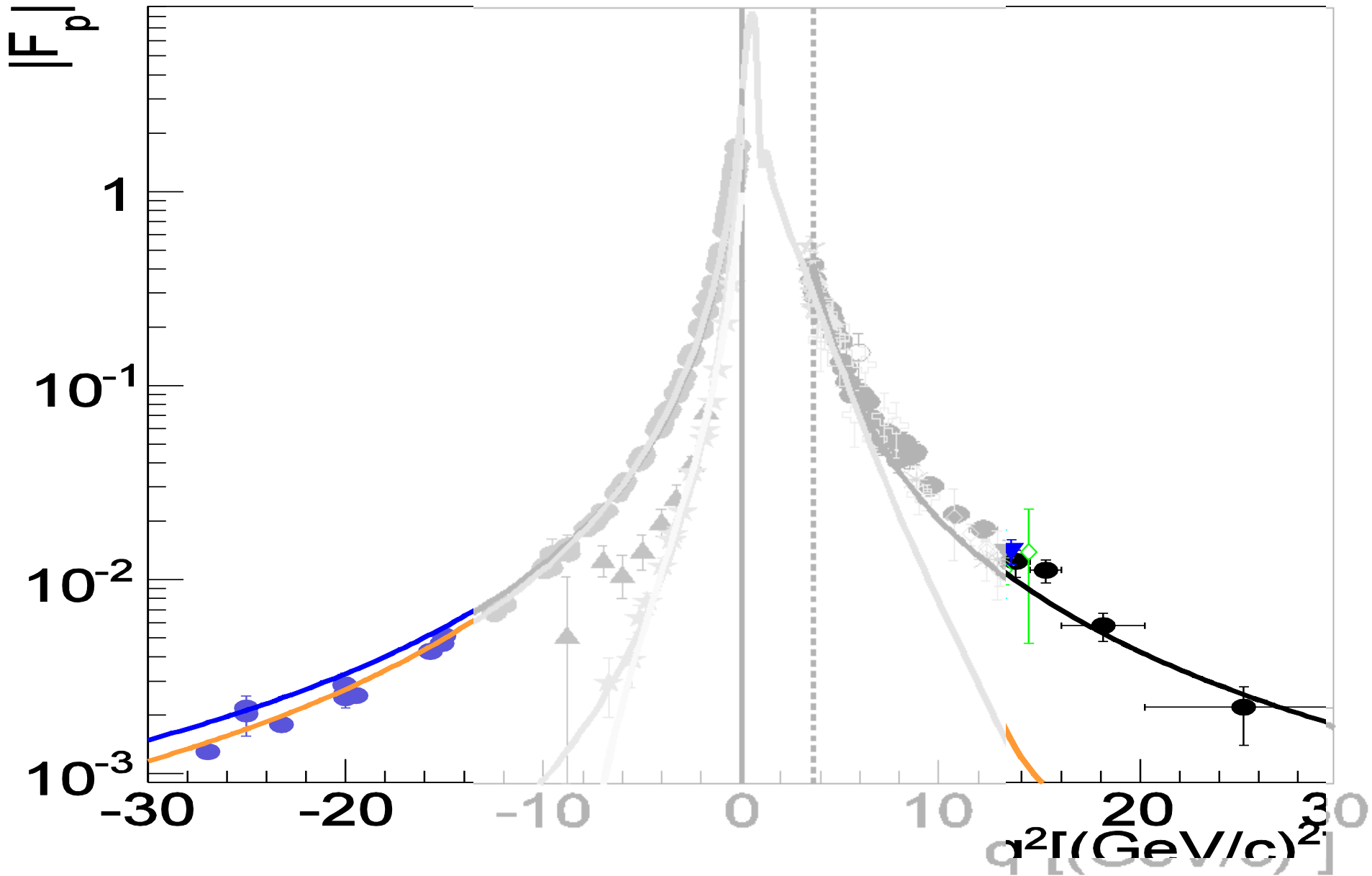
$Q(t)$: *time evolution of the charge distribution in the domain \mathcal{D} .*

Proton Form Factors



E.A. Kuraev, E. T-G, A. Dbeyssi, Phys.Lett. B712 (2012) 240

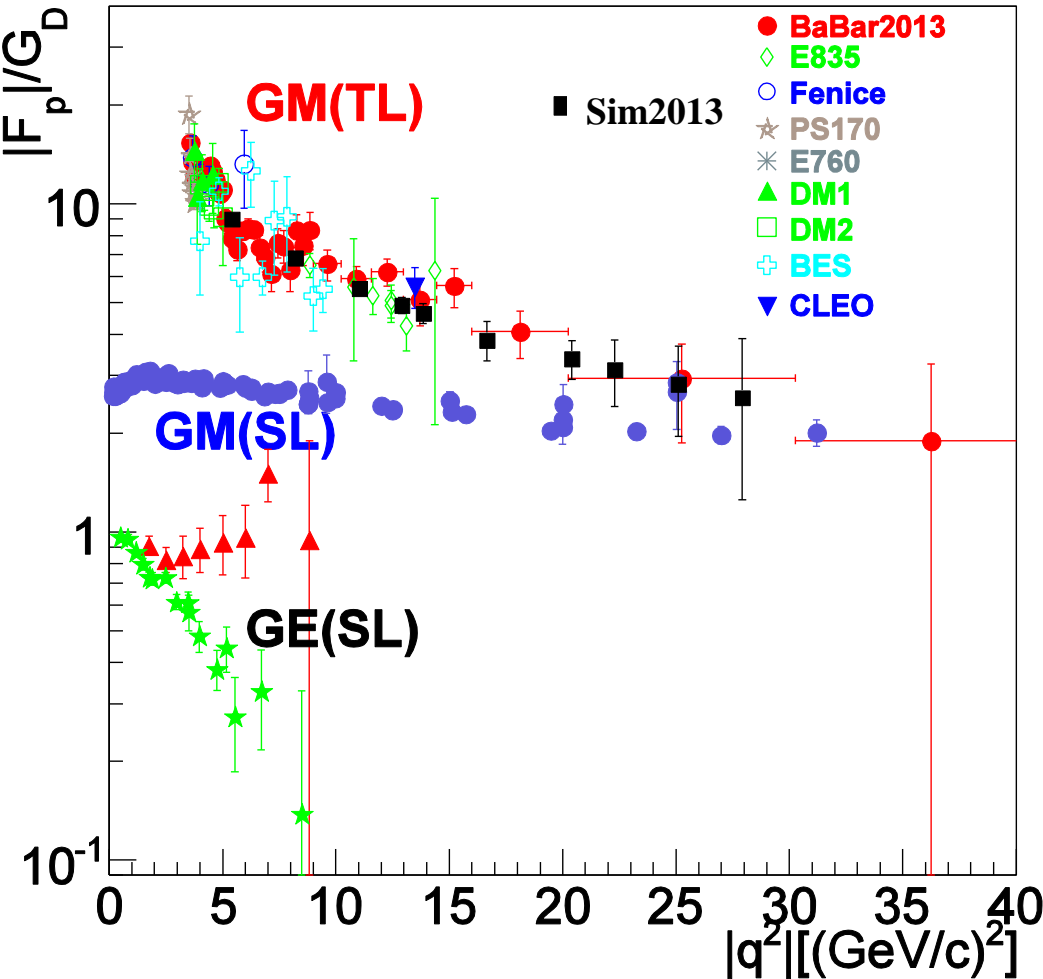
The asymptotic region



Proton form factors at large q^2

$\mathcal{L} = 2 \cdot 10^{32} \text{ cm}^{-2} \text{ s}^{-1}$
100 days

Phragmén-Lindelöf theorem



$$\lim_{q^2 \rightarrow -\infty} F^{(SL)}(q^2) = \lim_{q^2 \rightarrow \infty} F^{(TL)}(q^2)$$

space-like time-like
 $(e^- + p \rightarrow e^- + p)$ $(e^+ + e^- \leftrightarrow \bar{p} + p)$

– $F^{(TL)}(q^2) \rightarrow \text{real}$, if $q^2 \rightarrow \infty$

Applies to NN and $N\bar{N}$ Interaction
(Pomeranchuk theorem)
 $t=0$: not a QCD regime!

Analyticity
Connection with QCD asymptotics?

E. T-G. and M. P. Rekalo, Phys. Lett. B 504, 291 (2001)

Nucleon Form Factor Experiments

Hall	Exp#	Title	E_e	Q_{\max}^2
A	E12-07-108	Precision Measurement of the Proton Elastic Cross Section at High Q^2	6.6 8.8 11	17,5 (14)
A	E12-07-109	Large Acceptance Proton Form Factor Ratio Measurements at 13 and 15 (GeV/c) ² using Recoil Polarization Method	6.6 8.8 11	12(14)
A	E12-09-019	Precision Measurement of the Neutron Magnetic Form Factor up to $Q^2 = 18.0$ (GeV/c) ² by the Ratio Method	4.4 6.6 8.8 11	13.5 (18)
A	E12-09-016	Measurement of the Neutron Electromagnetic Form Factor Ratio G_E^N / G_M^N at High Q^2	4.4 6.6 8.8	10.2
B	E12-07-104	Measurement of the Neutron Magnetic Form Factor at High Q^2 Using the Ratio Method on Deuterium	11	14
C	E12-11-009	The Neutron Electric Form Factor at Q^2 up to 7 (GeV/c) ² from the Reaction $2H(e,e'n)1H$ via Recoil Polarimetry	4.4 6.6 11	7



Patrizia Rossi

ECT* Trento – February 18-22, 2013

9



Conclusions



Jefferson Lab

VEPP-3

Novosibirsk

IHEP

BES



• Large activity both in Space and Time-like regions

• Unified models in SL and TL regions:

- describe proton, neutron, electric, magnetic FFs
- pointlike behavior at threshold?
- understand $GE, GM(SL) < GE, GM(TL)$;

• To measure

- zero crossing of GE/GM in SL? 2γ ? Proton radius?
- GE and GM separately in TL (PANDA)
- complex FFs in TL region: polarization



*Σας ευχαριστώ
για την προσοχή σας*

The polarization method (exp: 2000)

Transferred polarization is:

*C. Perdrisat et al,
JLab-GEp collaboration*

$$P_n = 0$$

$$\pm h P_t = \mp h 2\sqrt{\tau(1+\tau)} G_E^p G_M^p \tan\left(\frac{\theta_e}{2}\right) / I_0$$

$$\pm h P_l = \pm h (E_e + E_{e'}) (G_M^p)^2 \sqrt{\tau(1+\tau)} \tan^2\left(\frac{\theta_e}{2}\right) / M / I_0$$

Where, $h = |h|$ is the beam helicity

$$I_0 = (G_E^p(Q^2))^2 + \frac{\tau}{\epsilon} (G_M^p(Q^2))^2$$

$$\Rightarrow \frac{G_E^p}{G_M^p} = -\frac{P_t}{P_l} \frac{E_e + E_{e'}}{2M} \tan\left(\frac{\theta_e}{2}\right)$$

The simultaneous measurement of P_t and P_l reduces the systematic errors

PID and kinematical Cuts

A. DBEYSSI ,PhD 2013

s [GeV ²]	5.4	8.2	13.9
Total PID prob.	>99%	>99%	>99.9%
Individual PID _{<i>i</i>} prob.	>5%	>5%	>6%
Number of fired crystals	>5	>5	>5
$(\theta + \theta')$ [CMS]	[178°-182°]	[178°-182°]	[175°-185°]
$ \phi - \phi' $	[178°-182°]	[178°-182°]	[175°-185°]
Invariant mass [GeV]	No cut	> 2.14 GeV	> 2.5 GeV
Background [Events]	0	0	0

- *PID --> probability for the detected particle to be identified as the signal.*
- *PID information are taken from **EMC**, **STT**, **DIRC** and **MVD** subdetectors.*